# REVISITED ISOPERIMETRIC INEQUALITIES FOR THE p-CAPACITY AND APPLICATION TO THE MUSKAT PROBLEM 

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## ABSTRACT

We give a unified form to various isoperimetric inequalities of $p$-capacity type and we present an application to a Muskat problem with prescribed flux.

## 1. INTRODUCTION

Let $1<p<\infty$ be a real number, $N$ a positive integer and $p^{\prime}$ the conjugate of $p: \frac{1}{p}+\frac{1}{p^{\prime}}=1$. Let $\omega_{1} \subset \subset \omega_{2}$ be given bounded open sets in $\mathbb{R}^{N}$ having respective boundaries $\partial \omega_{1}=\gamma_{1}, \partial \omega_{2}=\gamma_{2}$ and Lebesgue measures $m_{1}, m_{2}$. We define the domain $\Omega=\omega_{2} \backslash \bar{\omega}_{1}$. We denote by $\xi_{1} \cdot \xi_{2}$ the inner product of $\xi_{1}$ and $\xi_{2} \in \mathbb{R}^{N}$, by $|\xi|$ the Euclidean norm of $\xi \in \mathbb{R}^{N}$, by $|\omega|$ the Lebesgue measure of a measurable subset $\omega \subset \mathbb{R}^{N}$ and by $\beta_{N}$ the Lebesgue measure of the unit ball of $\mathbb{R}^{N}$.

Let $\alpha$ be a function of $L^{\infty}(\Omega)$, positive almost everywhere with $\frac{1}{\alpha}=\alpha^{-1} \in L^{\infty}(\Omega)$. Let $u \in W^{1, p}(\Omega)$ be such that
(1.1) $u_{y \gamma 1}=$ constant $>u_{y^{2}}=$ constant
and $\sigma \in L^{p^{\prime}}(\Omega)^{N}$ be a vector field which is divergence free (in the sense of distributions):
(1.2) $-\operatorname{div} \sigma=0$ in $\Omega$.

Furthemore, we assume that the pair $(u, \sigma)$ satisfies the inequality

$$
\text { (1.3) } \sigma \cdot \nabla u \geq \alpha|\nabla u|^{p} \text { a.e. in } \Omega
$$

We will provide some examples of vector fields $\sigma$ with their underlying functions $\alpha$.

In this paper, we show that $(u, \sigma)$ satisfies a general isoperimetric inequality which brings in a function $U$ of $W^{1, p}(\tilde{\Omega})$ verifying

$$
\left\{\begin{array}{l}
-d i v \Sigma=-d i v \tilde{\mathscr{A}} U=0 \text { in } \tilde{\Omega}, \\
U_{\mid \tilde{\gamma} 1}=\text { constant }>U_{\mid \tilde{\gamma} 2}=\text { constant },
\end{array}\right.
$$

with

- $\tilde{\Omega}=\overline{\omega_{2}} / \overline{\widetilde{\omega}_{1}}$, where $\overline{\omega_{1}}$ are the balls of $\mathbb{R}^{N}$ centered at the origin and having the same measures as $\omega_{i}$ and $\tilde{\gamma}_{i}=$ $\partial \widetilde{\omega}_{i}$ for $i=1$ or 2 ,
- $\sum(x)=(\tilde{\mathcal{A}} U)(x)=\tilde{\alpha}(x)|\nabla U(x)|^{p-2} \nabla U(x)$ where $\tilde{\alpha}$ is the spherical radially increasing rearrangement of $\alpha$ on $\tilde{\Omega}$. We will give later on a precise definition of this rearrangement introduced by A. Alvino and G. Trombetti [AlTr1, AlTr2].

Remark 1. The constant values $U_{\mid \hat{y}_{1}}$ and $U_{\mid \hat{\gamma}_{2}}$ are not necessarily the same as $u_{\gamma_{1}}$ and $u_{\gamma_{\gamma_{2}}}$.

Remark 2. Indeed, the condition (1.1) can be replaced by

$$
u_{q{ }^{1}}=\text { constant }<u_{q \gamma 2}=\text { constant }
$$

(set $u^{\prime}=-u$ and $\sigma^{\prime}=-\sigma$ ).
Our main tool is the theory of rearrangement of functions of Sobolev type, introduced by G. Talenti [Ta].

We will present some applications of this general result to various problems of Mathematical Physics such as the Muskat problem, a model arising from Oil Engineering.

## 2. EXAMPLES

We give in this section some examples of vector fields $\sigma$ with their corresponding functions $\alpha$.

Let $A: \Omega \rightarrow \mathbb{R}^{N \times N}$ be a matrix with measurable coefficients defined almost everywhere in $\Omega$ and
$g: \Omega \times \mathbb{R} \times \mathbb{R}^{N} \times \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}(x, \zeta, \xi, \eta) \rightarrow g(x, \zeta, \xi, \eta)$ a function defined for almost every $x$ in $\Omega$ and for any $(\zeta, \zeta, \eta) \in \mathbb{R} \times \mathbb{R}^{N} \times \mathbb{R}^{N}$.

We assume that the function $g$ and the matrix $A$ are related by the following hypothesis:
$\left\{\begin{array}{l}\text { There exists a functions } \alpha: \Omega \rightarrow \mathbb{R}, \alpha \in L^{\prime \prime}(\Omega), \alpha>0 \text { and } \frac{1}{\alpha}=\alpha^{-1} \in L^{\prime \prime}(\Omega)\end{array}\right.$ such that a.e. $x \in \Omega, \forall \zeta \in \mathbb{R}, \forall \xi \in \mathbb{R}^{N}, g(x, \zeta, \xi, A(x) \xi) \cdot \xi \geq \alpha(x) \mid \xi \xi^{p}$.

For $v \in W^{1, p}(\Omega)$, we denote by $\mathcal{A v}$ the function defined for almost every $x$ in $\Omega$ by

$$
\mathscr{A} v(x)=g(x, v(x), \nabla v(x), A(x) \nabla v(x)) .
$$

We assume that $u$ is a function of $W^{1, p}(\Omega)$ verifying

- $\mathcal{A} u \in\left(L^{p \prime}(\Omega)\right)^{N}$,
- $-\operatorname{div} \mathscr{A} u=0$ in $\Omega$ (in the sense of distributions),
- $u_{\gamma_{1}}=$ constant $>u_{\gamma_{2}}=$ constant.

The vector field $\sigma$ is then given by $\mathcal{A} u$. We precise hereafter some operators $g$ :

1. We consider $g(x, \zeta, \xi, \eta) \equiv g(x, \eta)$ satisfying the condition $g(x, \eta) \cdot \eta \geq \rho(x)|\eta|^{p}$ with $\rho(x)>0$. We choose the matrix $A(x) \equiv a(x)$ Id with $a(x)>0$. We get

$$
\begin{aligned}
g(x, \zeta, \xi, A(x) \xi) \cdot \xi & =g(x, a(x) \xi) \cdot \xi \\
& =\frac{1}{a(x)} g(x, a(x) \xi) \cdot a(x) \xi \geq \rho(x)(a(x))^{p-1}|\xi|^{p}
\end{aligned}
$$

We suppose that the function $\alpha$ defined by $\alpha(x)=\rho(x)(a(x))^{p-1}$ for almost every $x \in \Omega$, belongs to $L^{\infty}(\Omega)$ as well as $\alpha^{-1}$. The equation satisfied by $u$ is

$$
-\operatorname{div} g(x, a(x) \nabla u(x))=0 \operatorname{in} \Omega
$$

2. We choose a function $g(x, \zeta, \xi, \eta) \equiv \rho(x)(\eta \cdot \xi)^{\frac{p}{2}-1} \eta$ with $\rho(x)>0$ and a matrix $A$ such that a.e. $x \in \Omega, \forall \xi \in \mathbb{R}^{N}, A(x) \xi \cdot \xi \geq a(x)|\xi|^{2}$ with $a(x)>0$. Then, we have

$$
g(x, \zeta, \xi, A(x) \xi) \cdot \xi=\rho(x)(A(x) \xi \cdot \xi)^{\frac{p}{2}} \geq \rho(x)(a(x))^{\frac{p}{2}}|\xi|^{p}
$$

In this case, we assume that the function $\alpha(x)=\rho(x)(a(x))^{\frac{p}{2}}$ is in $L^{\infty}(\Omega)$ as well as its inverse. The equation satisfied by $u$ is

$$
-\operatorname{div}\left[\rho(x)(A(x) \nabla u(x) \cdot \nabla u(x))^{\frac{p}{2}-1} A(x) \nabla u(x)\right]=0 \text { in } \Omega .
$$

We recover the operator of $[\mathrm{Bo1}, \mathrm{Bo} 2, \mathrm{BoMos} 1, \mathrm{Bo}-$ Mos2, Mos].
3. Let $g(x, \zeta, \xi, \eta) \equiv g(x, \zeta, \xi) \cdot \xi \geq C|\xi|^{p}$ where $C$ is a real positive constant. Here $\alpha$ is a function defined by $\alpha(x)=C$ for $x \in \Omega$. The function $u$ verifies the equation

$$
-\operatorname{div} g(x, u(x), \nabla u(x))=0 \text { in } \Omega .
$$

This case has been studied by J.I. Diaz [Di].

## 3. ISOPERIMETRIC INEQUALITIES

First, we will prove a general isoperimetric inequality. We recall that the pair $(u, \sigma)$ is a solution of the problem denoted by ( $\mathcal{P}$ ):

$$
(P)\left\{\begin{array}{l}
u \in W^{1, p}(\Omega), \\
\left.u\right|_{\gamma 1}=\text { constant }>\left.u\right|_{\gamma 2}=\text { constant, }, \\
\sigma \in\left(L^{p \prime}(\Omega)\right)^{N}, \\
-\operatorname{div} \sigma=0 \text { in } \Omega, \\
\sigma \cdot \nabla u \geq \alpha|\nabla u|^{p} \text { a.e. in } \Omega
\end{array}\right.
$$

and the pair $(U, \Sigma)$ is the solution of $(\tilde{\mathcal{P}})$ :
$(\tilde{P})\left\{\begin{array}{l}U \in W^{1, p}(\tilde{\Omega}), \sum=\tilde{\mathcal{A}} U=\tilde{\alpha}|\nabla U|^{p-2} \nabla U \in\left(L^{p \prime}(\Omega)\right)^{N}, \\ U_{\mid \tilde{\gamma} 1}=\text { constant }>U_{\mid \tilde{\gamma} 2}=\text { constant }, \\ -\operatorname{div} \Sigma=0 \text { in } \tilde{\Omega} .\end{array}\right.$
The function $\tilde{\alpha}$ is defined on $\tilde{\Omega}$ by $\tilde{\alpha}(x)=\alpha^{*}\left(\beta_{N}|x|^{N}-\left|\omega_{1}\right|\right)$ where $\beta_{N}$ is the measure of the unit ball of $\mathbb{R}^{N},\left|\omega_{1}\right|$ and $|x|$ are respectively the Lebesgue measure on $\mathbb{R}^{N}$ of $\omega_{1}$ and the Euclidean norm of the vector $\overrightarrow{O x}, \alpha^{*}$ is the unidimensional increasing rearrangement of $\alpha$, defined on $\overline{\Omega^{*}}=[0,|\Omega|]$ by $\alpha^{*}(0)=$ ess inf $\alpha$, $\alpha^{*}(|\Omega|)=\operatorname{ess} \sup \alpha, \alpha^{*}(s)=\inf \{\theta \in \mathbb{R},|\alpha<\theta| \geq s\}$ for $s \in(0,|\Omega|)$, with $|\alpha<\theta|=$ Lebesgue measure of $\{x \in \Omega$, $\alpha(x)<\theta\}$. Similarly, we define $\alpha_{*}$ the unidimensional decreasing rearrangement of $\alpha$ on $\overline{\Omega^{*}}=[0,|\Omega|]$ by $\alpha_{*}(s)=\alpha^{*}(|\Omega|-s)$ for $s \in[0,|\Omega|]$.

Lemma 1. Under the conditions (1.1), (1.2) and (1.3), we have for almost every $x \in \Omega$

$$
u_{\mid \gamma_{1}} \geq u(x) \geq u_{\mid \gamma^{2}} \text { and } \int_{\Omega} \sigma \cdot \nabla u d x>0
$$

Proof: Let us show that for almost every $x \in \Omega$, we have $u_{y^{1}} \geq u(x)$. Since $\sigma \in\left(L^{p \prime}(\Omega)\right)^{N}$ and $-\operatorname{div} \sigma=0$ in $\Omega$ in the sense of distributions, we get

$$
\begin{equation*}
\int_{\Omega} \sigma \cdot \nabla w d x=0, \quad \forall w \in W_{0}^{1, p}(\Omega) . \tag{3.1}
\end{equation*}
$$

We take as test function $w_{1}=\left(u-u_{\gamma_{1}}\right)_{+}$, the positive part of $u-u_{\mid \gamma 1}$. We have by (3.1) and then by (1.3),

$$
\begin{aligned}
& 0=\int_{\Omega} \sigma \cdot \nabla\left(u-u_{\mid \gamma 1}\right)_{+} d x=\int_{u \gg u_{\mid Y 1}} \sigma \cdot \nabla u d x \\
& \geq \int_{u>u_{\mid x 1}} \alpha(x)|\nabla u|^{p} d x=\int_{\Omega} \alpha(x)\left|\nabla\left(u-u_{\mid y 1_{1}}\right)_{+}\right|^{p} d x .
\end{aligned}
$$

Since $\alpha>0$, we deduce that $w_{1}=\left(u-u_{\mid \gamma 1}\right)_{+}=$constant $=$ 0 in $\Omega$. It holds $u(x) \leq u_{\mid y 1}$ for almost every $x \in \Omega$. Similarly, using the test function $w_{2}=\left(u_{\mathrm{y} 2}-u\right)_{+}=\left(u-u_{\mathrm{y} 2}\right)_{-}$, we show that for almost every $x \in \Omega, u_{\gamma \mid} \leq u(x)$.

Let us show now that the quantity $\int_{\Omega} \sigma \cdot \nabla u d x$ is positive. Assume that $\int_{\Omega} \sigma \cdot \nabla u d x$ vanishes. Using the inequality (1.3), we deduce easily that $u$ is constant in each connected components of $\Omega$. This is in contradiction with $u_{\mid \gamma 1}>u_{\mid \gamma 2}$.

The general isoperimetric inequality is given in the following proposition.

Proposition Assume that $(u, \sigma)$ verifies ( $\mathcal{P}$ ). Let

$$
C(u, \sigma)=\int_{\Omega} \sigma \cdot \nabla u d x
$$

Then, for all real numbers $\theta, \theta^{\prime}$ satisfying $u_{\gamma_{2}} \leq \theta \leq \theta^{\prime} \leq u_{\gamma_{1} 1}$, the following isoperimetric inequality holds

$$
\begin{gathered}
\theta^{\prime}-\theta \leq N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N}\left(\frac{C(u, \sigma)}{u_{\mid y \mathrm{l}}-u_{\mid \gamma 2}}\right)^{\frac{p^{\prime}}{p}} \\
\int_{\mu\left(\theta^{\prime}\right)}^{\mu(\theta)}\left(s+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-\mu\left(\theta^{\prime}\right)\right) d s
\end{gathered}
$$

with $\mu(\theta)=|u>\theta|=$ measure of $\{x \in \Omega, u(x)>\theta\}$, $\left.\mu\left(\theta^{\prime}\right)=\left|u>\theta^{\prime}\right|, m_{1}=\left|\omega_{1}\right|\right\}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-\mu\left(\theta^{\prime}\right)\right)$ indicates the value of the function $\left(\alpha^{*}\right)^{\frac{-p^{*}}{p}}$ at the point $s-\mu\left(\theta^{\prime}\right)$.

Proof: We follow L. Boukrim [Bo1, Bo2] for the proof of this proposition. By Lemma 1, any value of $u$ is in $\left[u_{y_{2}}, u_{y_{11}}\right]$. What follows is valid for almost every $\tau \in\left(u_{\gamma 2}, u_{\gamma 1}\right)$. We set
$z_{\tau}=\tau-(u-\tau)_{-}=\left\{\begin{array}{l}u \text { if } u \leq \tau \\ \tau \text { if } u>\tau,\end{array}\right.$
$v_{\tau}=\frac{u-u_{\gamma 2}}{u_{\gamma \gamma 1}-u_{\gamma \gamma 2}}-\frac{z_{\tau}-u_{\gamma \gamma 2}}{\tau-u_{\gamma \gamma 2}}= \begin{cases}\left(u-u_{\mid \gamma 2}\right)\left(\frac{1}{u_{\mid \gamma 1}-u_{\gamma \gamma 2}}-\frac{1}{\tau-u_{\mid \gamma 2}}\right) & \text { if } u \leq \tau \\ \frac{u-u_{\mid \gamma 2}}{u_{\gamma \gamma 1}-u_{\gamma \gamma 2}}-1 & \text { if } u>\tau .\end{cases}$
Since $u \in W^{1, p}(\Omega)$, we have $v_{\tau} \in W_{0}^{1, p}(\Omega)$. Taking $w=$ $\nu_{\tau}$ in (3.1), it follows

$$
\begin{aligned}
& 0=\int_{\Omega} \sigma \cdot \nabla v_{\tau} d x=\left(\frac{1}{u_{\mid \gamma_{1}}-u_{\mid \gamma_{2}}}-\frac{1}{\tau-u_{\mid \gamma_{2}}}\right) \\
& \int_{u \leq \tau} \sigma \cdot \nabla u d x+\frac{1}{u_{\gamma_{1}}-u_{q_{\gamma_{2}}}} \int_{u \nu \tau} \sigma \cdot \nabla u d x,
\end{aligned}
$$

that is

$$
\int_{u \leq \tau} \sigma \cdot \nabla u d x=\frac{\tau-u_{y_{2}}}{u_{q_{1}}-u_{q_{\gamma 2}}} C(u, \sigma) .
$$

In consequence, one has

$$
\begin{equation*}
-\frac{d}{d \tau} \int_{u>\tau} \sigma \cdot \nabla u d x=\frac{d}{d \tau} \int_{u \leq \tau} \sigma \cdot \nabla u d x=\frac{C(u, \sigma)}{u_{\mid \gamma 1}-u_{\mid \gamma 2}} . \tag{3.2}
\end{equation*}
$$

On the other hand, for $h>0$

$$
\begin{aligned}
\frac{1}{h} \int_{\tau<u \leq \tau+h}|\nabla u| d x & =\frac{1}{h} \int_{\tau \ll u \tau+h} \alpha^{\frac{-1}{p}} \alpha^{\frac{1}{p}}|\nabla u| d x \\
& \leq\left[\frac{1}{h} \int_{\tau<u \leq \leq+h} \alpha^{\frac{-p^{\prime}}{p}} d x\right]^{\frac{1}{p^{\prime}}}\left[\left.\frac{1}{h} \int_{\tau<u \leq \tau+h} \alpha\right|^{2}| |^{p} d x\right]^{\frac{1}{p}} \\
& \leq\left[\frac{1}{h} \int_{\tau<u \leq \tau+h}^{\frac{-p^{\prime}}{p}} d x\right]^{\frac{1}{p^{\prime}}}\left[\frac{1}{h} \int_{\tau \ll \leq \tau+h} \sigma \cdot \nabla u d x\right]^{\frac{1}{p}}
\end{aligned}
$$

The first inequality above arises from the Hölder inequality and the second one comes from the condition (1.3). Letting $h$ tend to 0 , one gets at the limit
$-\frac{d}{d \tau} \int_{u>\tau}|\nabla u| d x \leq\left[-\frac{d}{d \tau} \int_{u>\tau} \alpha^{\frac{-p^{\prime}}{p}} d x\right]^{\frac{1}{p^{\prime}}}\left[-\frac{d}{d \tau} \int_{u \geqslant \tau} \sigma \cdot \nabla u d x\right]^{\frac{1}{p}}$.
Using the relation (3.2), we are led to

$$
\begin{equation*}
-\frac{d}{d \tau} \int_{u>\tau}|\nabla u| d x \leq\left[-\frac{d}{d \tau} \int_{u>\tau} \alpha^{\frac{-p^{\prime}}{p}} d x\right]^{\frac{1}{p^{p^{\prime}}}}\left(\frac{C(u, \sigma)}{u_{\mid \gamma 1}-u_{y 22}}\right)^{\frac{1}{p}} . \tag{3.3}
\end{equation*}
$$

Thanks to the isoperimetric inequality for the generalized perimeter of De Giorgi relative to $\Omega$ of the set $\{u\rangle$ $\tau\}$, denoted by $P_{\Omega}(u>\tau)$ [De] and a result of Fleming and Rishel [FIRi], we have with $\mu(\tau)=|u>\tau|$

$$
\begin{gathered}
-\frac{d}{d \tau} \int_{u>\tau}|\nabla u| d x=P_{\Omega}(u>\tau)= \\
=P_{R^{N}}\left(\{u>\tau\} \cup \bar{\omega}_{1}\right) \geq N \beta_{N^{\frac{1}{N}}}\left(\mu(\tau)+m_{1}\right)^{1-\frac{1}{N}}
\end{gathered}
$$

since the set $\{u>\tau\}$ does not meet $\gamma_{2}$ and its boundary includes $\gamma_{1}$. Therefore, by (3.3)

$$
N \beta_{N^{\frac{1}{N}}}\left(\mu(\tau)+m_{1}\right)^{1-\frac{1}{N}} \leq\left[-\frac{d}{d \tau} \int_{u \nu \tau} \alpha^{\frac{-p^{\prime}}{p}} d x\right]^{\frac{1}{p^{\prime}}}\left(\frac{C(u, \sigma)}{u_{y_{11}}-u_{y_{2}}}\right)^{\frac{1}{p}}
$$

Furthermore, thanks to the derivation formula (see Rakotoson and Temam [RaTe])

$$
\frac{d}{d \tau} \int_{\mu>\tau} \alpha^{\frac{-p^{\prime}}{p}} d x=\mathcal{W}^{\prime}(\mu(\tau)) \mu^{\prime}(\tau)
$$

where $\mathcal{W}^{\prime}(s)=\left(\alpha^{\frac{-p^{\prime}}{p}}\right)_{\tau_{u}}(s)$, that is, the relative rearrangement of $\alpha^{\frac{-p^{p}}{p}}$ with respect to $u$ defined by J. Mossino and R. Temam [MosTe], we obtain
$1 \leq N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N}\left(\frac{C(u, \sigma)}{u_{y \gamma 1}-u_{l y 2}}\right)^{\frac{p^{\prime}}{p}}\left(\mu(\tau)+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}} \mathcal{W}^{\prime}(\mu(\tau))\left(-\mu^{\prime}(\tau)\right)$.
Integrating the inequality (3.4) between $\theta$ and $\theta^{\prime}$, we get
$\theta^{\prime}-\theta \leq N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N}\left(\frac{C(u, \sigma)}{u_{y_{1}}-u_{y_{2}}}\right)^{\frac{p^{\prime}}{p}}$
$\int_{\theta}^{\theta^{\prime}}\left(\mu(\tau)+m_{1}\right)^{\frac{p^{\prime}}{N^{\prime}}-p^{\prime}} \mathcal{W}^{\prime}(\mu(\tau))\left(-\mu^{\prime}(\tau)\right) d \tau$
$\leq N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N}\left(\frac{C(u, \sigma)}{u_{\mid r 1}-\eta_{\gamma 2}}\right)^{\frac{p^{\prime}}{p}} \int_{0}^{[s]} x_{\left[\mu\left(\theta^{\prime}\right), \mu(\theta)\right]}(s)\left(s+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{p}}\left(\alpha^{\frac{-p^{\prime}}{p}}\right)_{*}(s) d s$.
According to a result of Rakotoson [Ra], the integral

$$
\int_{0}^{|s|} \chi_{\left[\mu\left(\theta^{\prime}\right), \mu(\theta)\right]}(s)\left(s+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{\frac{-p^{\prime}}{p}}\right)_{\%}(s) d s
$$

is bounded above by

$$
\begin{aligned}
& \int_{0}^{[s]}\left(\chi_{\left[\mu\left(\theta^{\prime}\right), \mu(\theta)\right]}(\cdot)\left(\cdot+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\right)_{*}(s)\left(\alpha^{\frac{-p^{\prime}}{p}}\right)_{*}(s) d s \\
& =\int_{0}^{|s|} \chi_{\left[0, \mu(\theta)-\mu\left(\theta^{\prime}\right)\right]}(s)\left(s+m_{1}+\mu\left(\theta^{\prime}\right)\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}(s) d s \\
& =\int_{\mu\left(\theta^{\prime}\right)}^{\mu(\theta)}\left(s+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-\mu\left(\theta^{\prime}\right)\right) d s .
\end{aligned}
$$

Consequently
$\theta^{\prime}-\theta \leq N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N}\left(\frac{C(u, \sigma)}{u_{y y^{2}}-u_{y 2}}\right)^{\frac{p^{\prime}}{p}} \int_{\mu\left(\theta^{\prime}\right)}^{\mu(\theta)}\left(s+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-\mu\left(\theta^{\prime}\right)\right) d s$,
and this ends the proof of the propostion.
Applying this proposition with $\theta=u_{\gamma_{2} 2}$ and $\theta^{\prime}=u_{\mid \gamma 1}$, one obtains the

Theorem 1. Assume that $(u, \sigma)$ and $(U, \Sigma)$ verify respectively $(\mathcal{P})$ and $(\tilde{\mathcal{P}})$. Set

$$
\begin{aligned}
& C(u, \sigma)=\int_{\Omega} \sigma \cdot \nabla u d x \text { and } \tilde{C}(U, \Sigma)= \\
& =\int_{\Omega} \Sigma \cdot \nabla U d x=\int_{\Omega} \tilde{\alpha}|\nabla U|^{p} d x=\tilde{C}(U) .
\end{aligned}
$$

Then one obtains the isoperimetric inequality

$$
\begin{gathered}
\frac{\left(u_{\gamma \gamma 1}-u_{\gamma_{2}}\right)^{p}}{C(u, \sigma)} \leq N^{-p} \beta_{N} \frac{-p}{N}\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}}= \\
=\frac{\left(U_{\mid \hat{r}_{1}}-U_{\mid \hat{r}_{2}}\right)^{p}}{\tilde{C}(U, \Sigma)}
\end{gathered}
$$

with $m_{i}=\left|\omega_{i}\right|(i=1,2)$.
Proof: Taking $\theta=u_{y 2}$ and $\theta^{\prime}=u_{\mid \gamma 1}$ in the Proposition, one gets

$$
\begin{gathered}
u_{\mid \gamma 1}-u_{\mid \gamma 2} \leq N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N}\left(\frac{C(u, \sigma)}{u_{\mid \gamma 1}-u_{\gamma_{2}}}\right)^{\frac{p^{\prime}}{p}} \\
\int_{\mu\left(y_{y>x}\right)}^{\mu\left(u_{\gamma_{r}}\right)}\left(s+m_{1}\right)^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-\mu\left(u_{\mid \gamma_{1}}\right)\right) d s .
\end{gathered}
$$

According to Lemma 1, the inequality $u(x) \leq u_{\mid y_{1}}$ holds for almost every $x \in \Omega$. In consequence, we have $\mu\left(u_{\gamma_{1} 1}\right)=0$ and $\mu\left(u_{\gamma_{2}}\right) \leq|\Omega|$. One gets the announced inequality by bounding above the previous integral between 0 and $\mu\left(u_{\gamma_{2}}\right)$ by the integral between 0 and $|\Omega|$. The equality in the theorem is classical.

## 4. SOME APPLICATIONS

We denote by ( $p, \alpha$ )-capacity of $\Omega=\omega_{2} \backslash \bar{\omega}_{1}$, the quantity $\int_{\Omega} \alpha|\nabla v|^{p} d x$ where $v$ is the solution of - div $\left(\alpha|\nabla v|^{p-2} \nabla v\right)=0$ in $\Omega, \quad v_{\mid y_{1}}=1$ and $v_{\mid \gamma 2}=0$. We give below some applications of Theorem 1 .

### 4.1. The $(p, \alpha)$-capacity problem

We assume that $u_{p_{i}}=C_{i}(i=1$ or 2$)$ are given constants with $C_{1}>C_{2}$. We consider the problem ( $\left.\tilde{q}\right)$ with the same constants $U_{\mid \bar{n}}=C_{i}$. We obtain the isoperimetric inequality

$$
\begin{aligned}
\int_{\Omega} \sigma \cdot \nabla u d x=C(u, \sigma) & \geq \tilde{C}(U, \Sigma)=\int_{\Omega} \tilde{\alpha}|\nabla U|^{p} d x \\
& =N^{p} \beta_{N^{N}} \frac{p}{N}\left(C_{1}-C_{2}\right)^{p}\left[\int_{m_{1}}^{m_{1}} s^{\frac{p^{\prime}}{N}-p^{p}}\left(\alpha^{*}\right)^{-\frac{p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{-p}{p^{\prime}}}
\end{aligned}
$$

With $\sigma=\alpha|\nabla u|^{p-2} \nabla u, C_{1}=1$ and $C_{2}=0$, we recover the isoperimetric inequality for the ( $p, \alpha$ )-capacity given in [ AlTr 2 ] and $[\mathrm{Fe}]$.

### 4.1.1. Application

Let $\alpha_{n}$ be a sequence of functions defined on $\Omega$ such that their unidimensional increasing rearrangement $\alpha_{n}^{*}$ satisfy

$$
\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha_{n}^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s \rightarrow 0
$$

when $n$ tends to infinity. If $\left(\sigma_{n}, u_{n}\right)$ verifies the problem $(\mathcal{P})$ of the $(p, \alpha)$-capacity, then we have

$$
\int_{\Omega} \sigma_{n} \cdot \nabla u_{n} d x \rightarrow \infty
$$

Let's precise this application by assuming for instance that $\alpha_{n}$ takes two values: $\alpha_{n}=A_{n}^{1}$ in $\Omega_{n}^{1}$ with $\left|\Omega_{n}^{1}\right|=\frac{1}{n}$ and $\alpha_{n}=A_{n}^{2}>A_{n}^{1}$ in $\Omega_{n}^{2}=\Omega \backslash \Omega_{n}^{1}$. Hence

$$
\begin{aligned}
\left(\frac{p^{\prime}}{N}-p^{\prime}+1\right) \int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha_{n}^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s & =\left(A_{n}^{1}\right)^{\frac{-p^{p}}{p}}\left[\left(m_{1}+\frac{1}{n}\right)^{\frac{p^{\frac{1}{N}}}{} p^{\prime}+1}-\left(m_{1}\right)^{\frac{p^{\prime}}{N^{\prime}} p^{p^{\prime}+1}}\right] \\
& +\left(A_{n}^{2}\right)^{\frac{-p}{p}}\left[\left(m_{2}\right)^{\frac{p^{\prime}}{N}-p^{\prime}+1}-\left(m_{1}+\frac{1}{n}\right)^{\frac{p^{\prime}}{N} p^{\prime}+1}\right]
\end{aligned}
$$

if $p^{\prime} \neq \frac{N}{N-1}$ and
$\int_{m_{1}}^{m_{2}} s^{-1}\left(\alpha_{n}^{*}\right)^{\frac{-p^{*}}{p}}\left(s-m_{1}\right) d s=\left(A_{n}^{1}\right)^{\frac{-p^{\prime}}{p}} \ln \left(1+\frac{1}{n m 1}\right)+\left(A_{n}^{2}\right)^{\frac{-p^{\prime}}{p}} \ln \left(\frac{m_{2}}{m_{1}+\frac{1}{n}}\right)$
if $p^{\prime}=\frac{N}{N-1}$. In consequence, taking the equivalents, we get

$$
\begin{gathered}
\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha_{n}^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s \cong \\
\cong m_{1}^{\frac{p^{\prime}}{N}-p^{\prime}} \frac{1}{n}\left(A_{n}^{1}\right)^{\frac{-p^{\prime}}{p}}+\frac{m_{2} \frac{p}{}^{\frac{p^{\prime}}{N}}-p^{\prime}+1}{} \frac{m_{1} \frac{p^{\prime}}{N^{\prime}+p^{\prime}}}{\frac{p^{\prime}}{N}-p^{\prime}+1}\left(A_{n}^{2}\right)^{\frac{-p^{\prime}}{p}}
\end{gathered}
$$

if $p^{\prime} \neq \frac{N}{N-1}$ and

$$
\int_{m_{1}}^{m_{2}} s^{-1}\left(\alpha_{n}^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s \cong \frac{1}{m_{1}} \frac{1}{n}\left(A_{n}^{1}\right)^{\frac{-p^{\prime}}{p}}+\ln \left(\frac{m_{2}}{m_{1}}\right)\left(A_{n}^{2}\right)^{\frac{-p^{\prime}}{p}}
$$

if $p^{\prime}=\frac{N}{N-1}$. In this case, in order to let the following integral

$$
\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha_{n}^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s
$$

tend to zero, it is enough to take $A_{n}^{2} \rightarrow \infty$ and $\frac{1}{n}\left(A_{n}^{1}\right)^{\frac{-p^{\prime}}{p}} \rightarrow 0$. We can choose for instance $A_{n}^{2} \rightarrow \infty$ whereas $A_{n}^{1} \rightarrow 0$ but with the condition $A_{n}^{1} \gg\left(\frac{1}{n}\right)^{p-1}$ (e.g. $A_{n}^{1}=\left(\frac{1}{n}\right)^{p-2}$ if $p \geq 2$ ).

### 4.2. The prescribed flux problem

We denote by $Q(u, \sigma)$ the quantity

$$
Q(u, \sigma)=\frac{C(u, \sigma)}{u_{\gamma \gamma 1}-u_{\gamma_{2}}}
$$

For regular open sets $\Omega$, for regular $u$ and suitable $\sigma$, the quantity $Q(u, \sigma)$ is a physical parameter (see the remark below). It is the total flux. For this reason, we also call «flux» the quantity $Q(u, \sigma)$ without any regularity assumption on $\Omega, u$ or $\sigma$. Assume that $u$ and $U$ satisfy ( $\mathcal{P}$ ) and $(\tilde{P})$ as well as the condition $\frac{C(u, \sigma)}{u_{y_{1}}-u_{y_{2}}}=Q=$ $=\frac{\tilde{C}(U, \Sigma)}{U_{\mid \tilde{\gamma_{1}}}-U_{\mid \overrightarrow{\gamma_{2}}}}$. The value $Q>0$ is given but the values of $u_{\gamma_{1}}, u_{\mid \gamma 2}, U_{\mid \tilde{y} 1}$ and $U_{\mid \tilde{\gamma} 2}$ remain undetermined. The Theorem 1 gives an optimal estimate for the variation of $u$, that is, a precise comparaison of the quantities $u_{\mid \gamma 1}-u_{\gamma^{2} 2}$ and $U_{\mid \tilde{\gamma} 1}-U_{\mid \vec{\gamma} 2}$ :
$u_{\mid \gamma 1}-u_{\mid \gamma 2} \leq Q^{\frac{p^{\prime}}{p}} N^{-p^{\prime}} \beta_{N} \frac{-p^{\prime}}{N} \int_{m_{1}}^{m_{1}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s=U_{\mid \tilde{p} 1}-U_{\mid \hat{\gamma} 2}$.
In particular, if $u_{y 2}=U_{\mid \gamma^{2}}=0$, one obtains an optimal estimate for $u_{\gamma 1}$.

### 4.2.1. Application

Let $\alpha_{n}$ be a sequence of functions defined on $\Omega$ such that their unidimensional increasing rearrangements $\alpha_{n}^{*}$ satisfy

$$
\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha_{n}^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s \rightarrow 0
$$

when $n$ tends to infinity. Let $\left(\sigma_{n}, u_{n}\right)$ be any solution of the prescribed flux problem $(\mathcal{P})$ and such that $\left(u_{n}\right)_{\mid \gamma^{2}}=0$. Then, one has

$$
\left(u_{n}\right)_{\mid r 1}=\underset{x \in \Omega}{\operatorname{ess} \sup }\left|u_{n}(x)\right| \rightarrow 0
$$

If there exists $\alpha$ such that for any $n \in \mathbb{N}$, one has $\alpha_{n}$ $\geq \alpha$, then

$$
\alpha \int_{\Omega}\left|\nabla u_{n}\right|^{p} d x \leq \int_{\Omega} \alpha_{n}\left|\nabla u_{n}\right|^{p} \leq \int_{\Omega} \sigma_{n} \cdot \nabla u_{n} d x=Q\left(u_{n}\right)_{\mid \gamma 1} \rightarrow 0
$$

and finally $u_{n} \rightarrow 0$ in $W^{l, p}(\Omega)$ and $L^{\infty}(\Omega)$ (strongly).
Remark 3. Let $n$ be the unitary outer normal to $\Omega$ at $\gamma_{1} \cup \gamma_{2}$. We assume that $\Omega, u, \sigma$ are regular enough in order to define

$$
Q^{\prime}(\sigma)=\int_{\gamma 1} \sigma \cdot n d \gamma
$$

(where $d \gamma$ is the measure on $\gamma_{1} \cup \gamma_{2}$ ) and in order to apply the Green formula. It appears that the quantity $Q^{\prime}(\sigma)$ is in fact $Q(u, \sigma)$. This equality is shown in

Lemma 2. We assume that $\Omega, u, \sigma$ are regular enough. Then we have

$$
C(u, \sigma)=\left(u_{y_{1} 1}-u_{q_{2}}\right) Q^{\prime}(\sigma)
$$

Proof: Let's remark that

$$
\int_{\gamma^{2}} \sigma \cdot n d \gamma=-\int_{\gamma 1} \sigma \cdot n d \gamma
$$

which is a straightforward consequence of the Green formula:

$$
0=\int_{\Omega} d i v \sigma d x=\int_{\gamma 1 \cup \gamma 2} \sigma \cdot n d \gamma
$$

This yields

$$
C(u, \sigma)=-\int_{\Omega} u d i v \sigma d x+\int_{\gamma \mathrm{luy2} 2} u \sigma \cdot n d \gamma=\int_{\gamma \mathrm{luv} 2} u \sigma \cdot n d \gamma .
$$

Afterwards, using the conditions on $u$ on the boundary of $\Omega$, we obtain

$$
\int_{\gamma 1 \cup \gamma^{2}} u \sigma \cdot n d \gamma=\left(u_{\mid \gamma 1}-u_{\gamma \gamma 2}\right) Q^{\prime}(\sigma) .
$$

This leads to the formula of Lemma 2.

### 4.3. The problem of domains with given ( $p, \alpha$ )-capacity

We are given the boundary $\gamma_{1}=\partial \omega_{1}$ of a regular open set $\omega_{1}$, a real $Q>0$ and $\alpha: \mathbb{R}^{N} \backslash \overline{\omega_{1}} \rightarrow \mathbb{R}^{+}$a measurable function which is bounded as well as its inverse. Assume furthermore that $\alpha$ is rearrangebale in the sense

$$
\forall s>0, \exists t \geq 0,|\alpha<t| \geq s
$$

(see e.g. B. Simon [Si]). We consider the sphere $\tilde{\gamma}_{1}=\partial \tilde{\omega}_{1}$. Thus we define the unidimensional rearrangement of $\alpha$ on $(0,+\infty)$ by
$\left.\alpha^{*}(s)=\inf \left\{t, \operatorname{mes}\left\{x \in \mathbb{R}^{N} \backslash \bar{\omega}_{1}, \alpha(x)<t\right\} \geq s\right\}, \quad \forall s \in\right] 0,+\infty[$
and $\tilde{\alpha}: \mathbb{R}^{N} \backslash \overline{\tilde{\omega}_{1}} \rightarrow \mathbb{R}^{+}$by $x \rightarrow \alpha^{*}\left(\beta_{N}|x|^{N}-\left|\omega_{1}\right|\right)$. There exists a unique sphere $\Gamma_{2}$ such that, denoting by $\boldsymbol{\sigma}$ the annulus with boundaries $\tilde{\gamma}_{1}$ and $\Gamma_{2}$, the $(p, \tilde{\alpha})$-capacity of $\boldsymbol{\delta}$ is equal to $Q$ : the measure $M_{2}$ of the ball bounded by $\Gamma_{2}$ is the unique solution of the equation

$$
N^{p} \beta_{N} \frac{p}{N}\left[\int_{m_{1}}^{M_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{1-p}=Q
$$

If $\omega_{2}$ is any domain containing strongly $\omega_{1}\left(\omega_{2} \supset \supset \omega_{1}\right)$ and such that the ( $p, \alpha$ )-capacity of $\Omega=\omega_{2} \backslash \frac{2}{\omega_{1}}$ is equal to $Q$, the Theorem 1 says that we have necessarily $m_{2} \geq$ $M_{2}$. Indeed, by Theorem 1, we have

$$
\frac{1}{Q} \leq N^{-p} \beta_{N^{\prime}} \frac{-p}{N}\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left[\left(\alpha_{12}\right)^{*}\right]^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}}
$$

Since
$\left.\left(\alpha_{\Omega \Omega}\right)^{*}(s)=\inf \{t, \operatorname{mes}\{x \in \Omega, \alpha(x)<t\} \geq s\}, \quad \forall s \in\right] 0,|\Omega|[$ and

$$
\operatorname{mes}\{x \in \Omega, \alpha(x)<t\} \leq \operatorname{mes}\left\{x \in \mathbb{R}^{N} \backslash \bar{\omega}_{1}, \alpha(x)<t\right\}
$$

we have, for $s \in[0,|\beta|]$,
$\{t, \operatorname{mes}\{x \in \Omega, \alpha(x)<t\} \geq s\} \subset\left\{t, \operatorname{mes}\left\{x \in \mathbb{R}^{N} \backslash \bar{\omega}_{1}, \alpha(x)<t\right\} \geq s\right\}$.
Hence

$$
\forall s \in[0,|\Omega|], \quad \alpha^{*}(s) \leq\left(\alpha_{\mid \Omega}\right)^{*}(s)
$$

In consequence, we have

$$
\frac{1}{Q}=N^{-p} \beta_{N} \frac{-p}{N}\left[\int_{m_{1}}^{M_{2}} s^{\frac{p^{\prime}}{N} p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{\mathrm{l}}\right) d s\right]^{\frac{p}{p^{\prime}}}
$$

and

$$
\begin{aligned}
\frac{1}{Q} & \leq N^{-p} \beta_{N} \frac{-p}{N}\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left[\left(\alpha_{\mid \Omega}\right)^{*}\right]^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}} \\
& \leq N^{-p} \beta_{N} \frac{-p}{N}\left[\int_{m_{1}}^{m_{2}} s^{\frac{-p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}}
\end{aligned}
$$

Finally, we get

$$
\int_{m_{1}}^{M_{2}} s^{\frac{p^{*}}{N}-p^{*}}\left(\alpha^{*}\right)^{\frac{-p^{*}}{p}}\left(s-m_{1}\right) d s \leq \int_{m_{1}}^{m_{2}} s^{\frac{p^{*}}{N} p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{*}}{p}}\left(s-m_{1}\right) d s
$$

and therefore $M_{2} \leq m_{2}$ necessarily.

## 5. APPLICATION TO A MUSKAT PROBLEM

### 5.1. Recall of the Muskat problem and of previous results

Let $1<p<\infty, a_{1}, a_{2}$ and k be positive real constants. Let $\omega_{1}, \omega_{2}$ with $\omega_{1} \subset \subset \omega_{2}$ be bounded open sets of $\mathbb{R}^{N}$ with regular boundaries $\partial \omega_{1}=\gamma_{1}, \partial \omega_{2}=\gamma_{2}$. At time $t=0$ (initial time), we are given $\omega=\omega(0)$, an open set with regular boundary $\partial \omega(0)=\gamma(0)$ such that $\omega_{1} \subset \subset \omega(0) \subset \subset$ $\omega_{2}$. We set $\Omega=\omega_{2} \backslash \overline{\omega_{1}}, \Omega_{1}(0)=\omega(0) \backslash \overline{\omega_{1}}$ and $\Omega_{2}(0)=\omega_{2}$ $\backslash \overline{\omega(0)}$. From this initial position $\omega(0), \omega$ evolves with the time $t$ (we write $\omega=\omega(t)$ for this dependance on time $t$ ): its boundary $\partial \omega(t)=\gamma(t)$ moves according to the normal velocity

## (Q) ${ }_{1}$

$$
\begin{aligned}
v_{v}(x, t) & =-k a_{1}\left|\nabla u_{1}(x, t)\right|^{p-2} \nabla u_{1}(x, t) \cdot v(x, t)=-k a_{1}\left|\nabla u_{1}(x, t)\right|^{p-2} \frac{\partial u_{1}}{\partial v}(x, t) \\
& =-k a_{2}\left|\nabla u_{2}(x, t)\right|^{p-2} \nabla u_{2}(x, t) \cdot v(x, t)=-k a_{2}\left|\nabla u_{2}(x, t)\right|^{p-2} \frac{\partial u_{2}}{\partial \nu}(x, t)
\end{aligned}
$$

where $u_{1}(., t)=u_{1}(t)$ and $u_{2}(., t)$ are defined respectively in $\Omega_{1}(t)=\omega(\mathrm{t}) \backslash \overline{\omega_{1}}$ and in $\Omega_{2}(t)=\omega_{2} \backslash \overline{\omega(t)}$ and are the solutions of the following equations:
$(Q)_{2}$

$$
\begin{cases}-\operatorname{div}\left(a_{1}\left|\nabla u_{1}(t)\right|^{p-2} \nabla u_{1}(t)\right)=0 & \text { in } \Omega_{1}(t) \\ -\operatorname{div}\left(a_{2}\left|\nabla u_{2}(t)\right|^{p-2} \nabla u_{2}(t)\right)=0 & \text { in } \Omega_{2}(t) \\ u_{1}(t)=1 & \text { on } \gamma_{1} \\ u_{2}(t)=0 & \text { on } \gamma_{2} \\ u_{1}(t)=u_{2}(t) & \text { on } \gamma(t) \\ a_{1}\left|\nabla u_{1}(t)\right|^{p-2} \frac{\partial u_{1}}{\partial \nu}(t)=a_{2}\left|\nabla u_{2}(t)\right|^{p-2} \frac{\partial u_{2}}{\partial \nu}(t) & \text { on } \gamma(t) .\end{cases}
$$

The notation $\frac{\partial}{\partial \nu}$ stands for $\nabla_{x} \cdot v(x, t)$ where $v(x, t)$ is the unitary outer normal to $\Omega_{1}(t)$ at $\gamma(t)$. We denote by (Q) the problem $\left(Q_{1}\right)+\left(Q_{2}\right)$. For $p=2$, this problem is called «Muskat problem» [EO, Mu]. It models the mining of oil (fluid 2 occupying $\Omega_{2}(t)$ ) by injection of viscous water (fluid 1 occupying $\Omega_{1}(t)$ ). The model suggested above (with $1<p<\infty$ ) is a natural generalization of the Muskat one. When $p=2$ and $N>1, \mathrm{~F}$. Abergel and J. Mossino [AbMos] have proved the existence of regular solutions locally in time of the Muskat problem by means of the method of «normal variations». For $p=2$ and $N=2, \mathrm{~F}$. Yi [Y] has also given a proof of the existence and uniqueness of classical solution locally in time by the Newton iteration method. It is well known (see [Be]) that the stability or unstability of the interface $\gamma(t)$ corresponds respectively to the condition on the mobility ratio $M=\frac{a_{1}}{a_{2}}<1$ or $>1$. A mathematical interpretation of this stability condition can be found in [ $\mathrm{Ab}, \mathrm{AbMos}$ ]. When $a_{1}<a_{2}$ (and $1<p<\infty$ ), L. Boukrim and J. Mossino [BoMos1, BoMos2] have given isoperimetric inequalities by comparison with an evolution problem with spherical symmetry defined from (Q). They have given an optimal estimate of the «critical time», that is, the time after which no regular solution may exist, as well as an optimal estimate of the respective volumes of the domains $\Omega_{1}(t)$ and $\Omega_{2}(t)$.

### 5.2. A MUSKAT PROBLEM WITH PRESCRIBED FLUX

### 5.2.1. The most general statement

Let $p, \omega_{1}, \omega_{2}$ and $\omega(0)$ as in section 5.1. Let $\beta$ and $\xi$ be functions defined on $\mathrm{R}^{+}$such that $\xi$ is positive and $\beta$ does not vanish and is of constant sign. We assume that $\beta \xi$ $\in L^{1}\left(\mathbb{R}^{+}\right)$.

Let $\delta$ and $\eta$ be positive functions defined on ( $m_{1}, m_{2}$ ) with $m_{\mathrm{i}}=$ measure of $\omega_{\mathrm{i}}$, $(i=1,2)$, such that $\frac{1}{\delta}$ and $\frac{1}{\delta \eta} \in L^{1}\left(m_{1}, m_{2}\right)$. Let $a_{1}, a_{2}: \Omega \times\left[m_{1}, m_{2}\right] \rightarrow \mathbb{R}^{+}$be positive functions such that for any $m \in\left[m_{1}, m_{2}\right], a_{i}(., m)$ and $\frac{1}{a_{i}(., m)}$ are in $L^{\infty}(\Omega)$ and

$$
\begin{equation*}
\underset{x \in \Omega}{\operatorname{ess} \sup } a_{1}(x, m) \leq \underset{x \in \Omega}{\operatorname{ess} \inf } a_{2}(x, m) . \tag{5.1}
\end{equation*}
$$

From the initial position $\omega(0)$, $\omega$ evolces with the time $t$ : we write $\omega=\omega(t)$. We denote $\Omega_{1}(t)=\omega(t)$ \} $\overline{\omega_{1}}, \Omega_{2}(t)=\omega_{2} \backslash \overline{\omega(t)}$ and $m(t)=|\omega(t)|=$ measure of $\omega(t)$ and we define

$$
\begin{equation*}
\alpha(x, t)=\alpha_{i}(x, t)=a_{i}(x, m(t)) \text { if } x \in \Omega_{i}(t), i=1,2 \tag{5.2}
\end{equation*}
$$

We consider the very general problem such that at any time $t$, there exists a domain $\omega(t)$ with $\omega_{1} \subset \subset \omega(t) \subset \subset \omega_{2}$ and a pair $(u(t), \sigma(t)) \in W^{1, p}(\Omega) \times\left(L^{p^{\prime}}(\Omega)\right)^{N}$ satisfying
(5.3) $\quad-\operatorname{div} \sigma(t)=0, \sigma(t) \cdot \nabla u(t) \geq \alpha(t)|\nabla u(t)|^{p}$ in $\Omega$,

$$
\begin{equation*}
u(t)=0 \text { on } \gamma_{2}, \tag{5.4}
\end{equation*}
$$

$$
\begin{equation*}
\int_{\Omega} \sigma(t) \cdot \nabla u(t) d x=\xi(t) \eta(m(t)) u(t)_{|r|}, \frac{d}{d t} m(t)=(\beta \xi)(t)(\delta \eta)(m(t)) . \tag{5.6}
\end{equation*}
$$

This general problem is actually a statement coming from the problem that we set below and that we call «Muskat problem with prescribed flux»: $\partial \omega(t)=\gamma(t)$ moves with the normal speed

$$
\begin{aligned}
\left(\mathcal{R}_{1}\right) v_{v}(x, t) & =-\beta(t) \delta(m(t)) \sigma_{1}(x, t) \cdot v(x, t)= \\
& =-\beta(t) \delta(m(t)) \sigma_{2}(x, t) \cdot v(x, t)
\end{aligned}
$$

where $v(t)=v(x, t)$ is the unitary outer normal to $\Omega_{1}(t)$ at $\gamma(t)$ and $\sigma_{i}(., t)=\sigma_{i}(t)$ satisfy (accordingly with $u_{i}(t)$ )
$\left(\mathcal{R}_{2}\right)\left\{\begin{array}{l}-\operatorname{div} \sigma_{i}(t)=0, \sigma_{i}(t) \cdot \nabla u_{i}(t) \geq \alpha_{i}(t)\left|\nabla u_{i}(t)\right|^{p} \text { in } \Omega_{i}(t), i=1,2, \\ u_{1}(t)=\text { unknown constant on } \gamma_{1}, \\ u_{2}(t)=0 \text { on } \gamma_{2}, \\ u_{1}(t)=u_{2}(t) \text { and } \sigma_{1}(t) \cdot v(t)=\sigma_{2}(t) \cdot v(t) \text { on } \gamma(t), \\ \int_{r 1} \sigma_{1}(t) \cdot n d \gamma=\xi(t) \eta(m(t)) .\end{array}\right.$
Here and in the following, we denote by $n$ the normal to $\Omega$ at $\gamma_{1} \cup \gamma_{2}$. The last condition of $\left(\mathbb{R}_{2}\right)$ involves the flux $\int_{\gamma_{1}} \sigma_{1}(t) \cdot n d \gamma$ : imposing this flux is related classically to the fact that $u_{1}(t)$ is an undetermined constant on $\gamma_{1}$.

Let us prove that for any regular solution of $\left(\mathcal{R}_{1}\right)+\left(\mathcal{R}_{2}\right)$ , we have (5.6) and $\sigma(t)$ is divergence free in the sense of distributions. (By regular solution, we mean that the Green formula is valid.)

1. Proof of $-\operatorname{div} \sigma(t)=0$ in $\mathcal{D}(\Omega)=$ space of distributions on $\Omega$. For any test function $\varphi \in \mathcal{D}(\Omega)=C_{0}^{\infty}(\Omega)$,
$\int_{\Omega} \sigma(t) \cdot \nabla \varphi d x=\int_{\Omega_{1}} \sigma_{1}(t) \cdot \nabla \varphi d x+\int_{\Omega_{2}} \sigma_{2}(t) \cdot \nabla \varphi d x=$
$=\sum_{i=1}^{2} \int_{\Omega_{i}}\left(-\operatorname{div} \sigma_{i}(t)\right) \varphi d x+\int_{\gamma(t)} \varphi\left(\sigma_{1}(t) \cdot v(t)-\sigma_{2}(t) \cdot v(t)\right) d \gamma=0$.
2. Proof of (5.6)
2.a.

$$
\int_{\Omega} \sigma(t) \cdot \nabla u(t) d x=\sum_{i=1}^{2} \int_{\Omega_{1}}\left(-\operatorname{div} \sigma_{i}(t)\right) u_{i}(t) d x
$$

$$
\begin{aligned}
& +\int_{\gamma_{t}(t)} u_{1}(t)\left(\sigma_{1}(t) \cdot v(t)-\sigma_{2}(t) \cdot v(t)\right) d \gamma+u_{1}(t)_{\mid r 1} \int_{\gamma_{1}} \sigma_{1}(t) \cdot n d \gamma \\
& =u_{1}(t)_{\mid y 1} \int_{\gamma \gamma_{1}} \sigma_{1}(t) \cdot n d \gamma=\xi(t) \eta(m(t)) u_{1}(t)_{\mid \gamma r} .
\end{aligned}
$$

2.b.

$$
\begin{aligned}
\frac{d m}{d t}(t) & =\int_{\gamma(t)} v_{v}(x, t) d \gamma=-\beta(t) \delta(m(t)) \int_{\gamma(t)} \sigma_{1}(t) \cdot v(t) d \gamma \\
& =\beta(t) \delta(m(t)) \int_{\gamma_{1}} \sigma_{1}(t) \cdot n d \gamma=(\beta \xi)(t)(\delta \eta)(m(t))
\end{aligned}
$$

(The penultimate equality comes from the integration on $\Omega_{1}(t)$ of $-\operatorname{div} \sigma_{1}(t)=0$.)

### 5.2.2. Isoperimetric inequalities

We denote by ( $\mathcal{F}$ ) the statements (5.2.) to (5.6.). Our isoperimetric inequalities arise from the following theorem as a corollary of Theorem 1

Theorem 2 Let $p^{\prime}$ verifying $\frac{1}{p}+\frac{1}{p^{\prime}}=1$ and $\beta_{N}$ The measure of the unit ball of $\mathbb{R}^{N}$. Denote by $t_{c}$ the critical time of problem $(\mathcal{F})$ that is the maximal time of existence of solution to (F). Set

$$
\begin{gathered}
\Phi(m)= \\
=\left[\int_{m_{1}}^{m} \frac{p^{p^{\prime}}-p^{\prime}}{}\left(a_{1} * \frac{-p^{\prime}}{p}\left(s-m_{1}, m\right) d s+\int_{m}^{m_{2}} s^{\frac{p^{\prime}}{N^{\prime}}-p^{\prime}}\left(a_{2} *\right)^{\frac{-p^{\prime}}{p}}(s-m, m) d s\right]^{\frac{p}{p^{\prime}}} .\right.
\end{gathered}
$$

Then, for any time t for which there exists a solution to $(\mathcal{F})$ on $[0, t]$ (hence for any $t \leq t_{c}$ ), one has

$$
N^{p} \beta_{N} \frac{p}{N} \int_{0}^{t}\left(\Gamma_{1}(\tau)\right)^{p-1}|\beta(\tau)| d \tau \leq \operatorname{sgn}(\beta)\left[\int_{m(0)}^{m(t)} \frac{\Phi}{\delta}(s) d s\right]
$$

where $\Gamma_{1}(\tau)=u_{1}(\tau)_{\mid \gamma 1}$ and $\operatorname{sgn}(\beta)$ is the sign of $\beta$.
Proof: Let $t$ be a time for which there exists a solution to $(\mathcal{F})$ on $[0, t]$ and let $\tau \in[0, t]$. By (5.2) we have $\alpha(x, \tau)$ $=a_{i}(x, m(\tau))$ for $x \in \Omega_{i}(\tau)$ and we set

$$
C(u(\tau), \sigma(\tau))=\int_{\Omega} \sigma(x, \tau) \cdot \nabla u(x, \tau) d x, \quad \Gamma_{1}(\tau)=u_{1}(\tau)_{\mid \gamma 1}
$$

Then Theorem 1 ensures that

$$
\begin{equation*}
\frac{\left(\Gamma_{1}(\tau)\right)^{p}}{C(u(\tau), \sigma(\tau))} \leq N^{-p} \beta_{N} \frac{-p}{N}\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N} p^{\prime}}\left(\alpha^{*}(\tau)\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}} \tag{5.7}
\end{equation*}
$$

where $\alpha^{*}(\tau)$ is the increasing rearrangement of $\alpha(., \tau)$. It yields from (5.7) and (5.6) that
$N^{p} \beta_{N^{N}} \frac{p}{( }\left(\Gamma_{1}(\tau)\right)^{p-1} \leq \xi(\tau) \eta(m(\tau))\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}} p^{\prime}\left(\alpha^{*}(\tau)\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}}$.
By (5.6),

$$
\frac{d m(\tau)}{d \tau}=\beta(\tau) \delta(m(\tau)) \xi(\tau) \eta(m(\tau))
$$

and the previous inequality becomes

$$
\begin{gathered}
N^{p} \beta_{N^{N}}^{\frac{p}{N}}\left(\Gamma_{1}(\tau)\right)^{p-1} \leq \frac{1}{\beta(\tau) \delta(m(\tau))} \frac{d m(\tau)}{d \tau}\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N^{\prime}} p^{\prime}}\left(\alpha^{*}(\tau)\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}} \\
=\frac{\operatorname{sgn}(\beta)}{|\beta(\tau)| \delta(m(\tau))} \frac{d m(\tau)}{d \tau}\left[\int_{m_{1}}^{m_{1}} \frac{p}{}_{s^{\prime}}^{s^{p^{\prime}}}\left(\alpha^{*}(\tau)\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}}
\end{gathered}
$$

It follows from (5.1) that

$$
\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N} \cdot p^{\prime}}\left(\alpha^{*}(\tau)\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}} \leq \Phi(m(\tau))
$$

We get finally

$$
N^{p} \beta_{N^{\frac{p}{N}}}\left(\Gamma_{1}(\tau)\right)^{p-1}|\beta(\tau)| \leq \operatorname{sgn}(\beta) \frac{\Phi}{\delta}(m(\tau)) \frac{d m(\tau)}{d \tau}
$$

which, after integration between 0 and $t$, gives the inequality of the theorem above.

We will estimate the unknown value of $u(t)$ on $\gamma_{1}$, denoted $\Gamma_{1}(t)$, and show that this estimate is optimal by comparison of the problem ( $\mathcal{F}$ ) with an evolution problem $(\tilde{\mathcal{F}})$ with spherical symmetry defined from $(\mathcal{F})$. Clearly, we obtain $(\tilde{\mathcal{F}})$ by replacing $\omega_{1}, \omega=\omega(0)$ and $\omega_{2}$ respectively by $\tilde{\omega}_{i}, \tilde{\omega}=\tilde{\omega}(0)$ and $\tilde{\omega}_{2}$, the balls of $\mathbb{R}^{N}$ centered at the origin and having the same measures as $\omega_{1}, \omega(0)$ and $\omega_{2}$. We also replace $a_{i}(., m)$ by its radially increasing rearrangement $\tilde{a}_{i}(., m)$ on $\tilde{\Omega}=\tilde{\omega}_{2} \backslash \overline{\tilde{\omega}}_{1}$. We denote by $\tilde{\omega}(t)$, the ball related to $(\tilde{\mathcal{F}})$ at time $t \geq 0$. A priori the domain $\tilde{\omega}(t)$ is not the symmetrization of the domain $\omega(t)$ related to $(\mathcal{F})$. Its measure is denoted by $\tilde{m}(t)$. We set $\tilde{\Omega}_{1}(t)=$ $\tilde{\omega}(t) \backslash \overline{\tilde{\omega}_{1}}, \tilde{\Omega}_{2}(t)=\tilde{\omega}_{2} \backslash \overline{\tilde{\omega}(t)}$ and $\tilde{\alpha}(x, t)=\tilde{a}_{i}(x,|\tilde{\omega}(t)|)$ if $x$ $\in \tilde{\Omega}_{i}(t), i=1$, 2. Finally, in the statement of $(\mathcal{F}),(u(t)$, $\sigma(t))$ is replaced by the pair $(U(t), \Sigma(t))$ where $\Sigma(t)=\tilde{\alpha}(t)|\nabla U(t)|^{p-2} \nabla U(t) ;$ More precisely the problem $(\tilde{\mathcal{F}})$ is
$-\operatorname{div} \Sigma(t)=0$ in $\tilde{\Omega}$ with $\Sigma(t)=\tilde{\alpha}(t)|\nabla U(t)|^{p-2} \nabla U(t)$,

$$
\begin{equation*}
U(t)=0 \text { on } \tilde{\gamma}_{2}=\partial \tilde{\omega}_{2} \tag{5.9}
\end{equation*}
$$

$$
\begin{equation*}
U(t)=\text { undetermined constant on } \tilde{\gamma}_{1}=\partial \tilde{\omega}_{1} \tag{5.10}
\end{equation*}
$$

$$
\begin{align*}
& \int_{\tilde{\Omega}} \Sigma(t) \cdot \nabla U(t) d x=  \tag{5.11}\\
& =\xi(t) \eta(\tilde{m}(t)) U(t)_{\mid \tilde{\gamma},}, \frac{d}{d t} \tilde{m}(t)=(\beta \xi)(t)(\delta \eta)(\tilde{m}(t)) .
\end{align*}
$$

The problem $(\tilde{\mathcal{F}})$ admits a unique solution and we have the

Theorem 3. Let $\tilde{m}(t)$ be the measure of $\tilde{\omega}(t), \tilde{\Gamma}_{1}(t)$ be the undetermined value of $U(t)$ on $\tilde{\gamma}_{1}$ and $\tilde{t}_{c}$ the critical time for the symmetrized problem, that is, the time such that $\tilde{\gamma}(t)$ touches $\tilde{\gamma}_{1}($ if $\beta<0)$ or $\tilde{\gamma}_{2}$ (if $\beta>0$ ). The values $\tilde{t}_{c}, \tilde{m}(t)$ and $\tilde{\Gamma}_{1}(t)$ are explicitly given by
(1)

$$
\int_{0}^{\tilde{t}_{c}}(|\beta| \xi)(\tau) d \tau= \begin{cases}\int_{m_{1}}^{m_{1}(0)} \frac{d s}{(\delta \eta)(s)} & \text { if } \beta<0 \\ \int_{m(0)}^{m_{2}} \frac{d s}{(\delta \eta)(s)} & \text { if } \beta>0\end{cases}
$$

(2)

$$
\int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)}=\int_{0}^{t}(\beta \xi)(\tau) d \tau \text { for } 0 \leq t \leq \tilde{t}_{c}
$$

(3)

$$
N^{p} \beta_{N} \frac{p}{N}\left(\tilde{\Gamma}_{1}(t)\right)^{p-1}=(\Phi \eta)(\tilde{m}(t)) \xi(t) \text { for } 0 \leq t \leq \tilde{t}_{c} .
$$

(4) If ( $\mathcal{F}$ admits a solution on [0, t], then one has
(i) $t \leq \tilde{t}_{c}\left(\right.$ hence $\left.t_{c} \leq \tilde{t}_{c}\right)$,
(ii) $m(t)=\tilde{m}(t)$,
(iii) $\quad \Gamma_{1}(t) \leq \tilde{\Gamma}_{1}(t)$.

## Proof:

1. Let $\tilde{t}_{c}$ be the critical time for the symmetrized problem and $\tau \in\left[0, \tilde{t}_{c}\right]$. We have

$$
\frac{d \tilde{m}(\tau)}{d \tau}=(\beta \xi)(\tau)(\delta \eta)(\tilde{m}(\tau))
$$

hence

$$
\begin{equation*}
(\beta \xi)(\tau)=\operatorname{sgn}(\beta)(|\beta| \xi)(\tau)=\frac{1}{(\delta \eta)(\tilde{m}(\tau))} \frac{d \tilde{m}(\tau)}{d \tau} \tag{5.12}
\end{equation*}
$$

This leads to

$$
(|\beta| \xi)(\tau)=\frac{\operatorname{sgn}(\beta)}{(\delta \eta)(\tilde{m}(\tau))} \frac{d \tilde{m}(\tau)}{d \tau}
$$

that we integrate between 0 an $\tilde{t}_{c}$. We get the equality

$$
\begin{align*}
& \int_{0}^{\bar{c}_{c}}(|\beta| \xi)(\tau) d \tau=\operatorname{sgn}(\beta) \int_{0}^{\tau_{c}} \frac{1}{(\delta \eta)(\tilde{m}(\tau))} \frac{d \tilde{m}(\tau)}{d \tau} d \tau=  \tag{5.13}\\
& =\operatorname{sgn}(\beta) \int_{\tilde{m}(0)}^{\tilde{m}\left(\bar{z}_{\bar{c}}\right)} \frac{d s}{(\delta \eta)(s)}
\end{align*}
$$

We obtain the announced equality by using the definition of the critical time for the symmetrized problem and $m(0)=\tilde{m}(0)$.
2. Let $t \in\left[0, \tilde{t}_{c}\right]$. From (5.12), we obtain, since $m(0)$ $=\tilde{m}(0)$

$$
\int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)}=\operatorname{sgn}(\beta) \int_{0}^{t}(|\beta| \xi)(\tau) d \tau
$$

3. From Theorem 1 and (5.11), we obtain as in the proof of Theorem 2 for any $t \leq \tilde{t}_{c}$

$$
\begin{aligned}
N^{p} \beta_{N} \frac{p}{N}\left(\tilde{\Gamma}_{1}(t)\right)^{p-1} & =\xi(t) \eta(\tilde{m}(t))\left[\int_{m_{1}}^{m_{2}} s^{\frac{p^{\prime}}{N}-p^{\prime}}\left(\alpha^{*}\right)^{\frac{-p^{\prime}}{p}}\left(s-m_{1}\right) d s\right]^{\frac{p}{p^{\prime}}} \\
& =\xi(t) \eta(\tilde{m}(t)) \Phi((\tilde{m}(t)))
\end{aligned}
$$

4. We assume that the problem $(\mathcal{F})$ admits a solution on $[0, t]$.
(i) Let $\tau \in[0, t]$. By the second relation of (5.6)

$$
(\beta \xi)(\tau)=\operatorname{sgn}(\beta)(|\beta| \xi)(\tau)=\frac{1}{(\delta \eta)(m(\tau))} \frac{d m(\tau)}{d \tau}
$$

we have for any $t$ such that $0 \leq t \leq t_{c}$,

$$
\int_{0}^{t}(|\beta| \xi)(\tau) d \tau=\operatorname{sgn}(\beta) \int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)}
$$

Since $\beta(\tau)$ and $\frac{d m(\tau)}{d \tau}$ have the same signs, we are led to

$$
\int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)} \leq \int_{m(0)}^{m_{2}} \frac{d s}{(\delta \eta)(s)}=\int_{m(0)}^{\bar{m}\left(\tilde{c}_{c}\right)} \frac{d s}{(\delta \eta)(s)}
$$

if $\beta>0$ (the function $t \rightarrow \tilde{m}(t)$ is increasing). By the same way, we have

$$
\int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)} \geq \int_{m(0)}^{m_{1}} \frac{d s}{(\delta \eta)(s)}=\int_{m(0)}^{n_{n}^{n}\left(\tilde{c}_{c}\right)} \frac{d s}{(\delta \eta)(s)}
$$

if $\beta<0$. hence, using (5.13)

$$
\int_{0}^{t}(|\beta| \xi)(\tau) d \tau \leq \operatorname{sgn}(\beta) \int_{m^{m}(0)}^{m\left(\tilde{t}_{c}\right)} \frac{d s}{(\delta \eta)(s)}=\int_{0}^{\tilde{\tau}_{c}}(|\beta| \xi)(\tau) d \tau
$$

for any $t \leq t_{c}$. Particularly, we have for $t=t_{c}$,

$$
\int_{0}^{t_{c}}(|\beta| \xi)(\tau) d \tau \leq \int_{0}^{\tau_{c}}(|\beta| \xi)(\tau) d \tau
$$

That is to say that $t_{c} \leq \tilde{t}_{c}$ and consequently, there exists no regular solution after $\tilde{t}_{c}$.
(ii) We recall that for any $t$ such that $0 \leq t \leq \min$ $\left(t_{c}, \tilde{t}_{c}\right)=t_{c}$,

$$
\begin{aligned}
& \int_{\bar{m}(0)}^{\bar{m}(t)} \frac{d s}{(\delta \eta)(s)}=\operatorname{sgn}(\beta) \int_{0}^{t}(|\beta| \xi)(\tau) d \tau, \\
& \int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)}=\operatorname{sgn}(\beta) \int_{0}^{t}(|\beta| \xi)(\tau) d \tau .
\end{aligned}
$$

These two equalities and $m(0)=\tilde{m}(0)$ give

$$
\int_{\tilde{m}(0)}^{\tilde{m}(t)} \frac{d s}{(\delta \eta)(s)}=\int_{m(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)}=\int_{\tilde{m}(0)}^{m(t)} \frac{d s}{(\delta \eta)(s)}
$$

In conclusion, one has $\tilde{m}(t)=m(t)$ for any $t \leq t_{c}$. This proves that $\tilde{\omega}(t)$ is the symmetrized domain of $\omega(t)$ for any $t \leq t_{c}$.
(iii) For any $t \leq t_{c}$, we have

$$
\begin{aligned}
\left(\Gamma_{1}(t)\right)^{p-1} & \leq N^{-p} \beta_{N} \frac{-p}{N}(\Phi \delta)(m(t)) \xi(t) \\
\left(\tilde{\Gamma}_{1}(t)\right)^{p-1} & =N^{-p} \beta_{N}^{\frac{-p}{N}}(\Phi \delta)(\tilde{m}(t)) \xi(t) \\
& =N^{-p} \beta_{N} \frac{-p}{N}(\Phi \delta)(m(t)) \xi(t)
\end{aligned}
$$

because $\tilde{m}(t)=m(t)$ as it has just been proved. We deduce immediately that $\Gamma_{1}(t) \leq \tilde{\Gamma}_{1}(t)$ for any $t \leq t_{c}$.

Remark that from this inequality, one also has

$$
\begin{aligned}
C(u(t), \sigma(t)) & =\int_{\Omega} \sigma(t) \cdot \nabla u(t) d x=\Gamma_{1}(t) \eta(m(t)) \xi(t) \\
& \leq \tilde{\Gamma}_{1}(t) \eta(m(t)) \xi(t)=\tilde{C}(U(t), \Sigma(t))=\sum_{i=1}^{i=2} \int_{\tilde{\Omega}_{i}(t)} \tilde{a}_{i}\left|\nabla U_{i}(t)\right|^{p} d x
\end{aligned}
$$

for any $t \leq t_{c}$.

### 5.3. AN EXAMPLE OF MUSKAT PROBLEM WITH PRESCRIBED FLUX

### 5.3.1. The problem

Let $p, \omega_{i}, \omega(0), a_{i}(x, m), \beta, \xi, \delta$ and $\eta$ as in the section 5.2. For $i=1,2$ we consider functions $G_{i}$ veryfing
i) $\quad G_{i}:(x, m x) \in \Omega \times\left[m_{1}, m_{2}\right] \times \mathbb{R}^{N} \rightarrow G_{i}(x, m, \xi)$ $\in \mathbb{R}$ are Caratheodory functions (that is, measurable with respect to $x$, continuous with respect to ( $m, \xi$ )),
ii) For almost every $x \in \Omega$, for any $m \in\left[m_{1}, m_{2}\right], G_{i}$ ( $x, m$..) is strictly convex and admits a gradient $g_{i}(x, m,$.$) ,$
iii) There exists $c^{1}, c^{2}, c^{4}>0$ and $c_{i}^{3} \in L^{1}(\Omega)$ such that a.e. $x \in \Omega, \forall \xi \in \mathbb{R}^{N}$ and $\forall m \in\left[m_{1}, m_{2}\right]$

$$
\begin{gathered}
c^{4}|\xi|^{p} \leq G_{i}(x, m, \xi) \leq c^{2}|\xi|^{p}+c_{i}^{3}(x) \\
\left|g_{i}(x, m, \xi)\right| \leq c^{4}\left(1+|\xi|^{p-1}\right) \\
g_{i}(x, m, \xi) \cdot \xi \geq a_{i}(x, m)|\xi|^{p} .
\end{gathered}
$$

We set for $i=1,2$

$$
\sigma_{i}(x, t)=g_{i}\left(x, m(t), \nabla u_{i}(t)\right) .
$$

Then, we have for $u_{i}(t) \in W^{1, \mathrm{p}}\left(\Omega_{i}(\mathrm{t})\right)$

$$
\begin{aligned}
\sigma_{i}(x, t) \cdot \nabla u_{i}(t) & =g_{i}\left(x, m(t), \nabla u_{i}(t)\right) \cdot \nabla u_{i}(t) \\
& \geq a_{i}(x, m(t))\left|\nabla u_{i}(t)\right|^{p}=\alpha_{i}(t)\left|\nabla u_{i}(t)\right|^{p} .
\end{aligned}
$$

With such $\sigma_{\mathrm{i}}$, we consider the statements $\left(\mathcal{R}_{1}\right),\left(\mathcal{R}_{2}\right)$ of the Muskat problem with prescribed flux.

Example: We set

$$
G_{i}(x, m, \xi)=\frac{1}{p}\left(A_{i}(x, m) \xi \cdot \xi\right)^{\frac{p}{2}}
$$

where $A_{i}=\Omega \times\left[m_{1}, m_{2}\right] \rightarrow \mathbb{R}^{N \times N}$ are symmetric matrices such that for any $m \in\left[m_{1}, m_{2}\right], A_{i}(., m) \in \mathrm{L}^{\infty}(\Omega)^{N \times N}$ and

$$
\text { a.e. } x \in \Omega, \forall \xi \in \mathbb{R}^{N} A_{i}(x, m) \xi \cdot \xi \geq a_{i}(x, m)|\xi|^{2} .
$$

W recover the operator of [BoMos2].

### 5.3.2. Existence of solution for a given time $\boldsymbol{t}$

For a given time $t$, we obtain the problem ( $\mathbb{R}_{2}$ ), with $Q(t)=\xi(t) \eta(m(t))$
$\left(R_{2}\right)$

$$
\left\{\begin{array}{l}
-\operatorname{div} g_{i}\left(x, m(t), \nabla u_{i}(t)\right)=0 \text { in } \Omega_{i}(t), i=1,2 \\
u_{1}(t)=\text { unknown constant on } \gamma_{1}, \\
u_{2}(t)=0 \text { on } \gamma_{2}, \\
u_{1}(t)=u_{2}(t) \text { on } \gamma(t), \\
g_{1}\left(x, m(t), \nabla u_{1}(t)\right) \cdot v(t)=g_{2}\left(x, m(t), \nabla u_{2}(t)\right) \cdot v(t) \text { on } \gamma(t), \\
\int_{\gamma_{1}} g_{1}\left(x, m(t), \nabla u_{1}(t)\right) \cdot n d \gamma=Q(t) .
\end{array}\right.
$$

We relate to $\left(\mathbb{R}_{2}\right)$ the minimization problem $\left(\mathcal{M}_{t}\right)$ : $\operatorname{Inf}\{J(v), v \in V\}$ where
$V=\left\{v \in W^{1, p}(\Omega), v=0\right.$ on $\gamma_{2}, v_{\mid \gamma 1}=$ undetermined constant $\}$
and

$$
J(v)=\int_{\Omega_{1}(t)} G_{1}(x, m(t), \nabla v) d x+\int_{\Omega_{2}(t)} G_{2}(x, m(t), \nabla v) d x-Q(t) v_{\mid r 1} .
$$

Then $V$ is a closed subspace of the reflexive Banach space $W^{1, p}(\Omega)(1<p<\infty)$ with its usual norm. We equip $V$ with the «gradient» norm $\|v\|_{V}=\|\nabla v\|_{L^{\prime}(\Omega)}, v \in V$. By the Poincaré inequality, this norm is equivalent to the one induced by $W^{1, p}(\Omega)$. One cheks that $J$ is strictly convex, continuous and coercive on $V$. There exists, in consequence, a unique solution to the minimization problem $\left(\mathcal{M}_{l}\right)$. This solution, denoted by $u$, is characterized by the variational equation

$$
\left\{\begin{array}{l}
u \in V \\
\int_{\Omega_{4}(t)}^{u} g_{1}(x, m(t), \nabla u) \cdot \nabla v d x+\int_{\Omega_{2}(t)} g_{2}(x, m(t), \nabla u) \cdot \nabla v d x= \\
=\xi(t) \eta(m(t)) v_{\mid r \mathrm{r}}, \forall v \in V .
\end{array}\right.
$$

Finally, using the Green fromula, we get formally the problem ( $\mathcal{R}_{2}$ ).

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