

One-parameter family of cubic Kolmogorov system with an isochronous center

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ABSTRACT

We show the existence of an one-parameter family of cubic Kolmogorov system with an isochronous center in the realistic quadrant.

1. Introduction

A center is said to be isochronous if all periodic orbits in a neighborhood of it have the same period. Recently, several papers have been devoted to the study of conditions on polynomial systems so that its centers are isochronous [4], [3], [5].

The autonomous differential system on the plane given by

$$(1) \quad \begin{cases} \dot{x} = x F(x, y) \\ \dot{y} = y G(x, y) \end{cases}$$

known as Kolmogorov system, is frequently used to model the interaction of two species occupying the same ecological niche.

If F and G are linear (Lotka-Volterra-Gause model), then it is well known that there is at most one critical point in the interior of the realistic quadrant ($x \geq 0$,

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$y \geq 0$), and that there are no limit cycles. In particular, in this situation periodic solutions can only occur nested around a center.

For the classical Lotka-Volterra system, J. Waldvogel [7] proved the monotonicity of the period function and, in particular, the non-existence of isochronous centers. For the general case of quadratic Kolmogorov System we show (by use of the classification given in [2]) the non-existence of isochronous centers (Theorem A).

In the case that F and G are both quadratic, that is, for cubic Kolmogorov system, one might think by analogy that the behavior within the first quadrant is similar to that of a quadratic Kolmogorov system. In this paper we show that this is not the case (Theorem B).

A technique to prove the isochronicity of some centers is the Urabe's Criteria [6]. To apply these criteria we must find two functions $\phi, \psi \in C^2$ such that the system (1) can be transformed, with the change of variables

$$x = \phi(u) \quad \text{and} \quad u = \psi(t),$$

into an equivalent equation of the form

$$(2) \quad \ddot{u} + g(u) = 0.$$

If $g'(0) = 1$, then all the solutions of (2) are periodic of period 2π if and only if

$$\frac{1}{2} X^2 = \int_0^u g(\xi) d\xi \iff u = X + F(X),$$

where $F(X)$ is an even function of X and $\text{sgn}(X) = \text{sgn}(u)$.

2. Main results

Theorem A

If F and G are linear, then the system (1) has no isochronous centers.

Proof. If the point (x_0, y_0) is a center of (1), then there exists a non-singular linear transformation $\Psi : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ with $\Psi^{-1}(x_0, y_0) = (0, 0)$, such that (1) is conjugate with a system of the form:

$$(3) \quad \begin{cases} \dot{u} = -v + a_{20}u^2 + a_{11}uv + a_{02}v^2 \\ \dot{v} = u + b_{20}u^2 + b_{11}uv + b_{02}v^2. \end{cases}$$

In [2] Loud proved that the system of the form (3) has an isochronous center at the origin only in four cases. Moreover, in [4] the first rational integrals are shown. It follows that for each case there exists at most one invariant straight line as level curve. As Ψ preserve the invariant straight lines of (1), we conclude that the point (x_0, y_0) cannot be an isochronous center. \square

Theorem B

Let us consider the one real parameter family of cubic Kolmogorov system

$$(4) \quad \begin{cases} \dot{x} = x (1 - a - 2 x (1 - a) - a (x^2 - y^2)) \\ \dot{y} = y (-(1 - a) + 2 y (1 - a) - a (x^2 - y^2)). \end{cases}$$

Then, for $a > 1$, the point $(1/2, 1/2)$ is an isochronous center of period $\frac{2\pi}{\sqrt{a-1}}$

Proof. The system (4) has a first rational integral

$$(5) \quad H_a(x, y) = \frac{(a - 1) (1 - 2(x + y)) + a (-x + y)^2}{a (a - 1) (-1 + 2(x + y) + a (-x - y + 1)^2)}.$$

For $c \in] - 1/a, -1/a + \epsilon[$, $0 < \epsilon \ll 1$, the level curves $H_a(x, y) = c$ are ellipse and for $c = -1/a$ we have $H_a^{-1}(-1/a) = \{(1/2, 1/2)\}$. It follows then that the singularity $(1/2, 1/2)$ is a center.

If we consider the conjugation

$$\Psi(X, Y) = \frac{1}{2} \left(-X - \sqrt{a-1} Y + 1, -X + \sqrt{a-1} Y + 1 \right),$$

then (4) is conjugate to the system

$$(6) \quad \begin{cases} \dot{X} = \sqrt{a-1} Y (-1 + 2 X - a X^2) \\ \dot{Y} = \sqrt{a-1} (X - X^2 + Y^2 - a X Y^2). \end{cases}$$

By rescaling the time

$$(7) \quad t = \frac{1}{\sqrt{a-1}} T,$$

the system (6) is transformed into the below system (in which we have replaced the capital letters T, X and Y by t, x and y , respectively).

$$(8) \quad \begin{cases} \dot{x} = y (-1 + 2 x - a x^2) \\ \dot{y} = x - x^2 + y^2 - a x y^2. \end{cases}$$

Firstly, we try to find an appropriate change of coordinates, to transform the system into a second order differential equation of the form (2). For this, consider $x = \phi(u)$ with $u = u(t)$, both of class C^2 . Replacing at the first equation of (8) we obtain

$$\phi'(u) \dot{u} = -y (1 - 2\phi(u) + a\phi^2(u))$$

where $\phi'(u)$ is the derivative of ϕ with respect to u and \dot{u} is the derivative of u with respect to the time t .

By the analyticity of the period function (see [4]), is sufficient to consider a small neighborhood \mathcal{V} of the origin such that $1 - 2\phi(u) + a\phi^2(u) > 0$.

Taken

$$y = \frac{\phi'(u) \dot{u}}{-1 + 2\phi(u) - a\phi^2(u)}$$

in the second equation of (8), we get

$$(1) \quad (-1 + 2\phi(u) - a\phi^2(u)) \phi'(u) \ddot{u} + [(-1 + 2\phi(u) - a\phi^2(u)) \phi''(u) - 3(1 - a\phi(u)) \phi'(u)^2] \dot{u}^2 - \phi(u) (1 - a\phi(u)) (-1 + 2\phi(u) - a\phi^2(u))^2 = 0.$$

Therefore, to obtain an equation of the form (2) we must consider the following initial value problem

$$\begin{aligned} (-1 + 2\phi(u) - a\phi^2(u)) \phi''(u) - 3(1 - a\phi(u)) \phi'(u)^2 &= 0 \\ \phi(0) = 0, \phi'(0) &= 1. \end{aligned}$$

Integrating the above equation over the neighborhood \mathcal{V} , we obtain

$$\begin{aligned} \phi'(u) &= \sqrt{(1 - 2\phi(u) + a\phi^2(u))^3} \\ \phi(0) &= 0. \end{aligned}$$

This initial value problem has the solution

$$(10) \quad \phi(u) = \frac{1}{a} \left(1 + \frac{au - u - 1}{\sqrt{1 + 2u + (1-a)u^2}} \right).$$

With this ϕ , equation (9) becomes in the form (2) with

$$(11) \quad g(u) = \frac{\phi(u) (1 - \phi(u))}{\sqrt{1 - 2\phi(u) + a\phi^2(u)}}.$$

Using (10) and (11), the equation

$$\frac{1}{2} X^2 = \int_0^u g(\xi) d\xi$$

transforms into $s \cdot z = 0$, where

$$\begin{aligned} s &= (-2u - u^2 + au^2 + 2X + 2uX - 2auX + aX^2) \\ z &= (-2u - u^2 + au^2 - 2X - 2uX + 2auX + aX^2). \end{aligned}$$

From $s = 0$, we obtain u in terms of X

$$u = X + F(X)$$

where

$$F(X) = \frac{1 - \sqrt{1 + (1 - a) X^2}}{a - 1}.$$

As $F(X)$ is even function, the Urabe's Criteria is satisfied and consequently for $a > 1$ the origin of the system (8) is an isochronous center. By the linear part of system (8) we conclude that the period is 2π . By (7) the origin of (6) is an isochronous center of period $2\pi/\sqrt{a-1}$ and, by the conjugation Ψ , the singularity $(1/2, 1/2)$ of the system (4) is an isochronous center of the same period. \square

Remark. By (5) it is easy to see that for $c \in]-1/a, 0[$, the level curves of $H_a(x, y) = c$ are the periodic orbits (ellipse) in the basin of the centre and, for $c = 0$, is the parabola $1 - 2(x + y) + a(x - y)^2 = 0$, the boundary of the basin.

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References

1. G.R. Fowles, "Analytic Mechanics", Holt, Rinerhart and Winston, 1977.
2. W. Loud, Behavior of the Period of Solutions of certain plane autonomous Systems near centers, *Contributions to Differential Equations* **31** (1964), 21–36
3. P. Mardešić, L. Moser-Jauslin and C. Rousseau, *Darboux linearization and isochronous centers with a rational first integral*, Preprint, Université de Bourgogne, 76, October, 1995.
4. P. Mardešić, C. Rousseau and B. Toni, Linearization of isochronous centers, *J. Differential Equations* **121** (1995), 67–108.
5. L. Mazzi and M. Sabatini, *Commutators and Linearizations of Isochronous Centres*, Preprint, Università degli Studi di Trento, 482, February, 1996.
6. M. Urabe, "Potential forces which yield periodic motions of a fixed period", *J. Math. Mech.* **10** (1961), 569–578.
7. J. Waldvogel, The Period in the Lotka-Volterra System is Monotonic, *J. Math. Anal. Appl.* **114** (1986), 178–184.