

Periodic solutions of a non coercive hamiltonian system

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ABSTRACT

In this paper, we study the existence of periodic solutions of hamiltonian systems:

$$\dot{x} = J H'(t, x)$$

where the hamiltonian H is non coercive of the type

$$H(t, r, p) = f(|p - Ar|) + h(t) \cdot (r, p).$$

1. Introduction

Let be given a relativistic particle with charge e and mass at rest m_0 submitted to a constant uniform magnetic field B and a uniform electric field $E(t)$, its movement is governed by the hamiltonian equations:

$$\dot{r} = \frac{\partial H}{\partial p}(t, r, p), \quad \dot{p} = -\frac{\partial H}{\partial r}(t, r, p)$$

where H is the particle energy given in terms of the time t , the position r and the impulsion p by the formula:

$$H(t, r, p) = c \left[m_0^2 c^2 + \left| p - \frac{e}{2c} B \wedge r \right|^2 \right]^{1/2} - e E(t) \cdot r$$

where c is the light speed.

This leads us to study the existence of periodic solutions of the hamiltonian systems:

$$(\mathcal{H}) \quad \dot{x} = JH'(t, x)$$

where the hamiltonian H is non coercive of the type

$$H(t, r, p) = f(|p - Ar|) + h(t) \cdot (r, p)$$

with A a matrix of order n , $f : \mathbb{R}_+ \rightarrow \mathbb{R}$ is a non decreasing convex continuously differentiable function, $h : \mathbb{R} \rightarrow \mathbb{R}^{2n}$ is a forcing term and J is the antisymmetric matrix:

$$J = \begin{pmatrix} 0 & I_{\mathbb{R}^n} \\ -I_{\mathbb{R}^n} & 0 \end{pmatrix}.$$

In this case the dual least action principle seems to provide the best results in the simplest way.

2. Autonomous case

In this section we take $h = 0$ and we assume that the matrix A is non symmetric and that the function f satisfies the assumptions:

- (1) $\forall t \in \mathbb{R}_+, 0 = f(0) < f(t)$ and $f'(0) = 0$;
- (2) $\exists k > 0, \exists c \geq 0 : \forall t \in \mathbb{R}_+, f(t) \leq \frac{k}{2}t^2 + c$;
- (3) $\exists \alpha > 0, \exists K \in \left] \frac{k}{2\tau}, +\infty \right[$ such that

$$\forall t \in [0, \alpha], \frac{k}{2}t^2 \leq f(t);$$

where

$$\tau = \sup \{ b \cdot (A^* - A)a; a^2 + b^2 = 1; a, b \in \mathbb{R}^n \}$$

and A^* is the adjoint of A . The corresponding hamiltonian H is then continuously differentiable and we obtain:

Theorem 1

For all $T \in \left] \frac{\pi}{K\tau}, \frac{2\pi}{k(1+\|A\|^2)} \right[$, the hamiltonian system (\mathcal{H}) has a non trivial periodic solution with minimal period T .

Proof. We proceed by proving successive lemmas.

We denote by f^* the Legendre transformation of f :

$$f^*(s) = \sup \{st - f(t); t \in \mathbb{R}\}.$$

From the assumptions (2) and (3) we deduce easily the following lemma:

Lemma 1

f^* satisfies

$$(4) \quad \forall t \in \mathbb{R}, f^*(t) \geq \frac{1}{2k} t^2 - c,$$

$$(5) \quad \exists r > 0; \forall |t| \leq r, f^*(t) \leq \frac{1}{2K} t^2.$$

It is easy to show that the function H is convex and its Legendre transformation H^* is given for $(s, q) \in \mathbb{R}^n \times \mathbb{R}^n$ by:

$$H^*(s, q) = \begin{cases} f^*(|q|) & \text{if } s + A^*q = 0 \\ +\infty & \text{elsewhere.} \end{cases}$$

Now, let $T \in]\frac{\pi}{K\tau}, \frac{2\pi}{k(1+\|A\|^2)}[$. We consider the functional Φ in the space E_0 defined by

$$E_0 = \left\{ (-A^*v, v); v \in L^2(0, T; \mathbb{R}^n), \int_0^T v(t) dt = 0 \right\},$$

$$\Phi(w) = \frac{1}{2} \int_0^T \langle Jw, \pi w \rangle dt + \int_0^T H^*(w) dt,$$

where πw is the primitive of w with mean value zero:

$$\frac{d}{dt}(\pi w) = w, \int_0^T (\pi w)(t) dt = 0.$$

Lemma 2

Let, for $w \in L^2(0, T; \mathbb{R}^{2n})$

$$g(w) = \int_0^T H^*(w) dt.$$

The subdifferential of $g|_{E_0}$ at a point $w \in E_0$, where g is finite, is given by:

$$\bar{\partial}g(w) = \{u \in L^2(0, T; \mathbb{R}^{2n}); \exists x \in E_0^\perp, u(t) + x(t) \in \partial H^*(u(t)) \text{ a.e}\}$$

where $\bar{\partial}$ designates the subdifferential in E_0 .

Proof. It is clear that

$$\bar{\partial}g(w) = \partial(g + \delta_{E_0})(w)$$

where

$$\delta_{E_0}(w) = \begin{cases} 0 & \text{if } w \in E_0 \\ +\infty & \text{elsewhere.} \end{cases}$$

We have $\delta_{E_0}^* = \delta_{E_0^\perp}$ and

$$E_0^\perp = \mathbb{R}^{2n} + \{(v, Av); v \in L^2(0, T; \mathbb{R}^n)\}.$$

Consequently, for $u \in L^2(0, T; \mathbb{R}^{2n})$, we have

$$(g^* \nabla \delta_{E_0}^*)(u) = \inf_{x \in \mathbb{R}^{2n}} \int_0^T H(u + x) dt.$$

So, by the assumption (2), there exist $\alpha, \beta > 0$ such that

$$0 \leq (g^* \nabla \delta_{E_0}^*)(u) \leq \alpha \|u\|_{L^2}^2 + \beta$$

and, since $g^* \nabla \delta_{E_0}^*$ is convex, then it is continuous. Noting $u = (r, p)$, we have

$$(g^* \nabla \delta_{E_0}^*)(u) = \inf_{\xi \in \mathbb{R}^n} \int_0^T f(|p - Ar + \xi|) dt.$$

Therefore, by assumptions (2), (3) and the convexity of f , the function $g^* \nabla \delta_{E_0}^*$ is exactly.

In the other hand, g and δ_{E_0} are convex, l.s.c. and proper. Then for all $w \in E_0$, where g is finite, we have

$$\partial(g + \delta_{E_0})(w) = \partial g(w) + \partial \delta_{E_0}(w).$$

The result follows then from the equalities:

$$\partial \delta_{E_0}(w) = E_0^\perp \text{ and } \partial g(w) = \{u \in L^2(0, T; \mathbb{R}^{2n}); u(t) \in \partial H^*(w(t)) \text{ a.e.}\} \quad \square$$

Lemma 3

The function Φ has a global minimum on E_0 :

$$\exists \bar{w} \in E_0, \quad \min_{E_0} \Phi = \Phi(\bar{w}).$$

Proof. By the inequality (4) we have, for all $w \in E_0$;

$$\int_0^T H^*(w) dt \geq \frac{1}{2k(1 + \|A\|^2)} \|w\|_{L^2}^2 - cT.$$

By application of the Wirtinger inequality, we obtain

$$\Phi(w) \geq \frac{1}{2} \left\{ \frac{1}{k[1 + \|A\|^2]} - \frac{T}{2\pi} \right\} \|w\|_{L^2}^2 - cT$$

and since $T < \frac{2\pi}{k(1 + \|A\|^2)}$ and the space E_0 is reflexive the minimum of Φ on E_0 is achieved. \square

Lemma 4

We have

$$\min_{E_0} \Phi < 0.$$

Proof. By definition, we have

$$T > \frac{\pi}{K\tau} = \frac{\pi}{K} \inf \left\{ \frac{a^2 + b^2}{b \cdot (A^* - A)a} ; b \cdot (A^* - A)a > 0 \right\},$$

so there exist $a, b \in \mathbb{R}^n$ such that

$$b \cdot (A^* - A)a > 0, \quad a^2 + b^2 \leq r^2 \quad \text{and} \quad T > \frac{\pi(a^2 + b^2)}{K b \cdot (A^* - A)a}.$$

Let

$$v(t) = a \cos\left(\frac{2\pi}{T}t\right) + b \sin\left(\frac{2\pi}{T}t\right), \quad w(t) = (-A^*v(t), v(t)),$$

we have by the inequality (5) and easy calculation

$$\int_0^T f^*(|v|) dt \leq \frac{1}{2K} \|v\|_{L^2}^2 = \frac{a^2 + b^2}{4K},$$

$$\int_0^T \langle Jw, \pi w \rangle dt = -\frac{T}{2\pi} b \cdot (A^* - A)a,$$

consequently

$$\Phi(w) \leq \frac{1}{4\pi} b \cdot (A^* - A)a \left[\frac{\pi(a^2 + b^2)}{K b \cdot (A^* - A)a} - T \right] < 0$$

and so $\inf_{E_0} \Phi < 0$.

Now, let w be a point of E_0 where the minimum of Φ is achieved, then we have

$$0 \in -J\pi w + \bar{\partial}g(w).$$

From Lemma 2, there exist $\xi \in \mathbb{R}^{2n}$ and $h \in L^2(0, T; \mathbb{R}^n)$ such that

$$J(\pi w)(t) + \xi + (h(t), Ah(t)) \in \partial H^*(w(t)) \text{ a.e.}$$

By setting $x = J\pi w + \xi$, we obtain by the Fenchel reciprocity $\dot{x} = JH'(x)$. It is clear that x is T -periodic and, by lemma 4, x is non constant.

It remains to show that T is the minimal period of x . Assume that x and then w are $\frac{T}{m}$ -periodic with $m \geq 2$. Let $y(t) = w(\frac{t}{m})$, then $(\pi y)(t) = m(\pi w)(\frac{t}{m})$. This point y belongs to E_0 and verifies

$$\begin{aligned} \Phi(y) &= \frac{m}{2} \int_0^T \langle Jw, \pi w \rangle dt + \int_0^T H * (w) dt \\ &\leq m \Phi(w). \end{aligned}$$

So, by lemma 4, we have $\Phi(y) < \inf_{E_0} \Phi$ which is a contradiction, then T is the minimal period of x . \square

3. Non autonomous case

Here we assume that f is not constant and there exist $k > 0$ and $a \geq 0$ such that:

$$(6) \quad \forall t \in \mathbb{R}_+, \quad 0 \leq f(t) \leq \frac{k}{2}t^2 + a \quad \text{and} \quad f'(0) = 0,$$

and the function h is continuous, periodic with minimal period $T > 0$ and mean value zero.

Theorem 2

For all $T \in]0, \frac{\pi}{k[1+\|A\|^2]}[$, the hamiltonian system (\mathcal{H}) has a periodic solution with minimal period T .

Proof. We proceed as in section 2, so we omit some details. We consider the functional Φ over the space E defined by

$$E = E_0 + h$$

$$\Phi(w) = \frac{1}{2} \int_0^T \langle Jw, \pi w \rangle dt + \int_0^T H^*(t, w) dt$$

where E_0 is defined as in section 2.

For $w = (-A^*v, v) + h \in E$, we have

$$\Phi(w) = \frac{1}{2} \int_0^T \langle Jw, \pi w \rangle dt + \int_0^T f^*(|v|) dt.$$

From the assumption (6), we have

$$\int_0^T f^*(|v|) dt \geq \frac{1}{2k} \|v\|_{L^2}^2 - aT,$$

it follows then, by Hölder inequality, that there exists a constant c such that

$$\Phi(w) \geq \frac{1}{4\pi} \left[\frac{\pi}{k[1 + \|A\|^2]} - T \right] \|w\|_{L^2}^2 - c,$$

and then the global minimum of Φ over E is achieved at a point w . Therefore, there exists $\xi \in \mathbb{R}^{2n}$ and $r \in L^2(0, T; \mathbb{R}^n)$ such that

$$0 \in -J\pi w + \xi + (r, Ar) + \partial H^*(t, w) \quad a.e.$$

By using the Fenchel reciprocity, it is clear that the function $x = J\pi w + \xi$ is a periodic solution of (\mathcal{H}) with minimal period T . \square

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