

ON THE ORTHOGONALITY OF SOME SYSTEMS  
OF POLYNOMIALS

by

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1. INTRODUCTION. Consider the BERNOULLI and EULER polynomials defined by means of

$$(1.1) \quad \frac{te^{tx}}{e^t - 1} = \sum_{n=0}^{\infty} B_n(x) \frac{t^n}{n!}$$

$$(1.2) \quad \frac{2e^{tx}}{e^t + 1} = \sum_{n=0}^{\infty} E_n(x) \frac{t^n}{n!}.$$

We have  $B_n(0) = B_n$ , the BERNOULLI numbers and  $2^n E_n(0) = C_n$  the tangent coefficients.

TOUCHARD [6] defined a set of polynomials  $\Omega_n(z)$  so that we have the symbolic orthogonality

$$(1.3) \quad B^m \Omega_n(B) = K_n \delta_{nm} \quad (0 \leq m \leq n),$$

where

$$K_n = \frac{(-1)^n [n!]^4}{2^n (2n+1) (1, 3, 5 \dots (2n-1))^2}.$$

It is understood in (1.3) that after expansion  $B^k$  is to be replaced by  $B_k$ .

TOUCHARD also indicated that (1.3) leads to actual orthogonality. He proved

$$(1.4) \quad -\frac{\pi i}{2} \int_{c-\infty i}^{c+\infty i} \frac{\Omega_n(z) \Omega_m(z)}{\sin^2 \pi z} dz = 2^n K_n \delta_{nm}$$

where  $-1 < c < 0$ .

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CARLITZ [2] generalized (1.3) and (1.4) in the following manner. Define

$$\beta_n = \beta_n(\lambda) = \frac{B_{n+1}(\lambda) - B_{n+1}}{(n+1)\lambda}$$

(obviously  $\beta_n(0) = B_n$ ) and consider the polynomials defined by

$$(1.5) \quad \Omega_n^{(\lambda)}(z) = (-2)^n (\lambda+1)_n \binom{2n}{n}^{-1} F_n^\lambda(1-\lambda+2z)$$

where  $F_n^\lambda(z)$  is the PASTERNAK [5] polynomial. He proved

$$(1.6) \quad \beta^r \Omega_n^{(\lambda)}(\beta) = K_n^{(\lambda)} \delta_{rn} \quad (0 \leq r \leq n)$$

where

$$K_n^{(\lambda)} = \frac{(-1)^n (n!)^2 (1+\lambda)_n (1-\lambda)_n}{(2n+1) 2^n \{1, 3, 5, \dots, (2n-1)\}^2}.$$

We have further

$$(1.7) \quad \frac{\sin \pi \lambda}{2i\lambda} \int_{c-\infty}^{c+\infty} \frac{\Omega_n^{(\lambda)}(z) \Omega_m^{(\lambda)}(z)}{\sin \pi z \sin \pi(z-\lambda)} dz = 2^n K_n^{(\lambda)} \delta_{nm}$$

where  $-1 < c < 0$ ,  $0 \leq \lambda < 1$ .

CARLITZ also proved

$$(1.8) \quad (\frac{1}{2}C)^m A_n(\frac{1}{2}C) = (-2)^{-n} S_{nm} \quad (0 \leq m \leq n)$$

where

$$\frac{(1+t)^z}{(1-t)^{z+1}} = \sum_{n=0}^{\infty} A_n(z) t^n.$$

In this note we shall give different proofs of (1.6), (1.7) and (1.8) making use of known facts in the theory of continued fractions. This method has the advantage that given the moments  $\beta_n(\lambda)$  the polynomials of (1.5) and the weight function are obtained in a natural way.

2. *Proof of (1.6).* Let  $\Psi(z)$  be the logarithmic derivation of the gamma function of  $z$ . We have the asymptotic formula [4, 106]

$$\Psi(z+h) \sim \log z - \sum_{n=1}^{\infty} (-1)^n \frac{B_n(h)}{nz^n}, \quad |\arg z| < \pi.$$

From this formula and familiar properties of the BERNOULLI polynomials we can show that

$$(2.1) \quad \frac{1}{h} \left\{ \Psi(z+1) - \Psi(z+1-h) \right\} \sim \sum_{n=0}^{\infty} \frac{\beta_n(h)}{z^{n+1}}.$$

This formula can be proved directly from the generating function of  $\beta_n(h)$ ,

$$\sum_{n=0}^{\infty} \beta_n(h) \frac{t^n}{n!} = \frac{1}{h} \frac{1 - e^{th}}{1 - e^t}.$$

Now consider the continued fraction [7, p. 372]

$$(2.2) \quad \frac{1}{2} \left\{ \Psi(z+k) - \Psi(z+1-k) \right\} = \frac{k - \frac{1}{2}}{z +} \frac{a_1}{z +} \frac{a_2}{z +} \dots$$

where

$$a_n = \frac{n^2(n+1-2k)(n-1+2k)}{4(2n-1)(2n+1)}.$$

From it follows, after some elementary transformation, that the continued fraction associated with the series in (2.1) is given by

(2.3)

$$\frac{1}{h} \left\{ \Psi(z+1) - \Psi(z+1-h) \right\} = \frac{2}{2z+1-h+} \frac{c_1}{2z+1-h+} \frac{c_2}{2z+1-h+} \dots$$

where

$$(2.4) \quad C_n = \frac{n^2(n^2-h^2)}{4n^2-1}.$$

Thus the  $\beta_n(h)$  are the moments of the J-fraction (2.3) which is uniformly convergent in every closed domain not including  $z = 0, -1, -2, \dots$

Consequently we have the symbolic orthogonality

$$(2.5) \quad \begin{aligned} \beta^r \Phi_n(\beta) &= 0 & (r = 0, 1, \dots, n-1), \\ &\neq 0 & (r = n), \end{aligned}$$

where  $\phi_n$  are the denominators of the convergents of the continued fraction (2.3). They satisfy the recurrence relation

$$\phi_{n+1}(z) = (2z+1-h)\phi_n(z) + c_n\phi_{n-1}(z).$$

If we compare this recurrence relation with that satisfied by the PASTERNAK polynomials

$$(n+1)(n+m+1)F_{n+1}^m(z) = -(2n+1)zF_n^m(z) + n(n-m)F_{n-1}^m(z)$$

we see that

$$\phi_n(z) = \Omega_n^{(A)}(z).$$

This completes the proof of (1.6).

3. *Proof of (1.7).* To prove (1.7) we begin by changing (2.3) into a positive definite  $J$ -fraction. We get

$$(3.1) \quad \frac{i}{2h} \left\{ \Psi\left(\frac{1}{2}zi + \frac{1}{2} + \frac{1}{2}h\right) - \Psi\left(\frac{1}{2}zi + \frac{1}{2} - \frac{1}{2}h\right) \right\} = \frac{1}{z-} \frac{c_1}{z-} \frac{c_2}{z-} \dots$$

The denominators of the convergents of (3.1) are orthogonal polynomials with respect to the distribution function  $\Psi(x)$  obtained by inverting the STIELTJES transform

$$\frac{i}{2h} \left\{ \Psi\left(\frac{1}{2}zi + \frac{1}{2} + \frac{1}{2}h\right) - \Psi\left(\frac{1}{2}zi + \frac{1}{2} - \frac{1}{2}h\right) \right\} = \int_{-\infty}^{\infty} \frac{d\Psi(u)}{z-u}.$$

In fact we know that [7, p. 253]

$$\Psi(u) = \lim_{y \rightarrow 0} \int_0^u \operatorname{Im} \frac{i}{h} \left\{ \Psi(ix - y + \frac{1}{2} - \frac{1}{2}h) - \Psi(ix - y + \frac{1}{2} + \frac{1}{2}h) \right\} dx.$$

If  $\phi(u)$  exists such that  $\Psi'(u) = \phi(u)$  then

$$\phi(u) = \lim_{y \rightarrow +0} \operatorname{Im} \frac{i}{h} \left\{ \Psi(ix - y + \frac{1}{2} + \frac{1}{2}h) - \Psi(ix - y + \frac{1}{2} - \frac{1}{2}h) \right\}.$$

To calculate  $\phi(x)$  we make use of the formula [4, p. 102]

$$\Psi(z+1) = -C + \sum_{n=1}^{\infty} \left( \frac{1}{n} - \frac{1}{n+2} \right)$$

where  $C$  is EULER'S constant. We get

$$\begin{aligned} \phi(x) &= \frac{1}{\pi h} \sum_{n=1}^{\infty} \left\{ \frac{2n-1-h}{(2n-h-1)^2+x^2} - \frac{2n-1+h}{(2n+h-1)^2+x^2} \right\} \\ &= \frac{\sin \pi h}{2h} \frac{1}{\cos \frac{\pi}{2}(h-ix) \cos \frac{\pi}{2}(h+ix)}. \end{aligned}$$

Replace  $ix$  by  $2x - h + 1$  we get the weight function associated with the polynomials  $\phi_n(x)$ , i.e.,

$$\rho(x) = \frac{\sin \pi h}{2h} \frac{1}{\sin \pi(x-h) \sin \pi x}.$$

From (1.2) we have

$$\frac{2}{e^t + 1} = \sum_{n=0}^{\infty} \frac{C_n}{2^n} \frac{t^n}{n!}.$$

We thus have asymptotic formula

$$2 \int_0^{\infty} \frac{e^{-xt}}{e^t + 1} dt = 2 \int_0^{\infty} \operatorname{sech} t e^{-t(2x+1)} dt \simeq \sum_{n=0}^{\infty} \frac{2^{-n} C_n}{x^{n+1}}.$$

If we recall the continued fraction expansion [7, p. 206]

$$(3.2) \quad \int_0^{\infty} \operatorname{sech}^k t e^{-zt} dt = \frac{1}{z + \frac{1 \cdot k}{z + \frac{2 \cdot (k+1)}{z + \frac{3 \cdot (k+2)}{z + \dots}}}}.$$

it follows that  $2^{-n} C_n$  are the moments associated with the orthogonal polynomials defined by means of

$$f_{n+1}(x) = (2x + 1) f_n(x) + n^2 f_{n-1}(x)$$

$$f_0(x) = 1, f_1(x) = x + \frac{1}{2}.$$

But  $f_n(x) = n! A_n(x)$ . We thus have (1.8).

In the same way as in the proof of (1.7) we can obtain the actual orthogonality satisfied by the polynomials  $A_n(z)$ . This has been essentially obtained by different methods in [1] and [3].

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